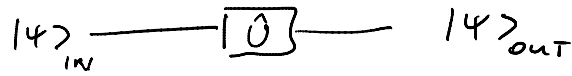


I.4. Single Qubit Gates

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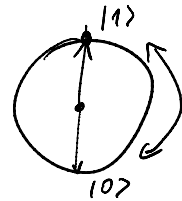
Circuit representation:



Work out how specific operations are represented on the Bloch sphere:

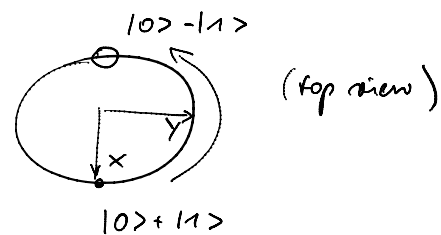
→ identity: $\hat{U} = \hat{I} = \mathbb{1}_2 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$

→ bit flip: $\hat{U} = \hat{X} = \sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \hat{=} \text{NOT operation}$
(rotation by π about x-axis)



→ phase flip: $\hat{U} = \hat{Z} = \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

$$\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \frac{1}{\sqrt{2}} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$



→ conjugate bit flip: $\hat{U} = \hat{Y} = \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$
(bit-phase flip)

$$|0\rangle \rightarrow -i|1\rangle$$

$$|1\rangle \rightarrow i|0\rangle$$

$\sigma_x, \sigma_y, \sigma_z$... Pauli matrices

$\{1, \sigma_x, \sigma_y, \sigma_z\}$ basis set for operators on a single qubit

Hadamard gate: generation of superposition from $|0\rangle$ or $|1\rangle$

$$\hat{H} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} = \frac{1}{\sqrt{2}} (\hat{X} + \hat{Z})$$

$\hat{=}$ rotation by π about $\begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix}$ -axis

$$|0\rangle \xrightarrow{\boxed{\hat{H}}} \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$$

$$|1\rangle \xrightarrow{\boxed{\hat{H}}} \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$$

\Rightarrow used in many quantum algorithms to prepare (initial) superposition state

Single Qubit Dynamics

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time-evolution given by Schrodinger equation

$$\boxed{\hat{H}|\psi\rangle = i\hbar \frac{\partial}{\partial t} |\psi\rangle}$$

e.g. single spin- $\frac{1}{2}$ particle in magnetic field:

$$\hat{H} = -\vec{\mu}_s \cdot \vec{B} = -\mu_s \frac{\hbar}{2} \begin{pmatrix} \hat{\sigma}_x \\ \hat{\sigma}_y \\ \hat{\sigma}_z \end{pmatrix} \cdot \begin{pmatrix} B_x \\ B_y \\ B_z \end{pmatrix} = -\frac{\mu_s \hbar}{2} (\hat{\sigma}_x B_x + \hat{\sigma}_y B_y + \hat{\sigma}_z B_z)$$

simplest case $\vec{B} = B_z \cdot \vec{e}_z$

$$\hat{H} = -\mu_s B_z \frac{\hat{\sigma}_z}{2}$$

energy eigenstates: $\hat{H}|0\rangle = E_0|0\rangle$

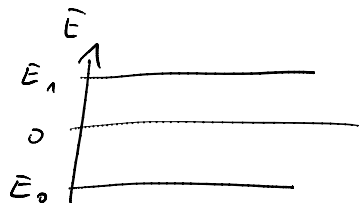
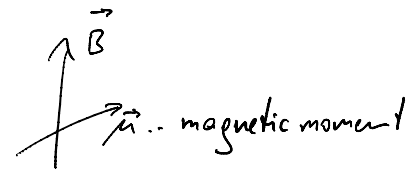
$$\hat{H}|1\rangle = E_1|1\rangle$$

$$E_0 = -\left(-\frac{g_s e \hbar}{2m_e}\right) B \cdot (-1) = -\frac{g_s \mu_B B}{2}$$

$$E_1 = \frac{g_s \mu_B B}{2}$$

$$\Delta E = E_1 - E_0 = g_s \mu_B B \hat{=} \hbar \omega_L$$

ω_L ... Larmor frequency



What's the evolution, if \vec{n} points in x-direction?

Larmor precession about z-axis!

$$|\psi(t)\rangle = e^{-\frac{i}{\hbar} \hat{H} t} |\psi(0)\rangle \quad \text{for time-indep. } \hat{H}$$

$$-\frac{i}{\hbar} \hat{H} = i \frac{2\mu_B B}{2\hbar} \vec{n} \cdot \vec{\sigma}$$

$$= i \frac{\omega}{2} \vec{n} \cdot \vec{\sigma}$$

$$\text{from } \vec{B} = |\vec{B}| \cdot \vec{n} = |\vec{B}| \cdot \begin{pmatrix} \sin \theta \cos \varphi \\ \sin \theta \sin \varphi \\ \cos \theta \end{pmatrix}$$

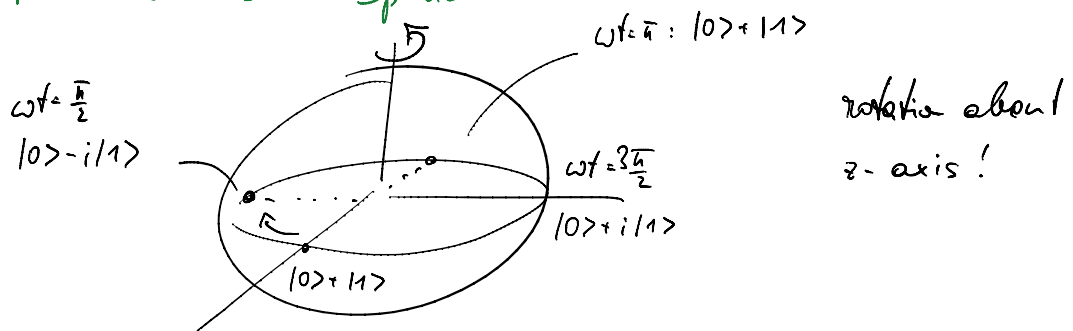
↑
unit vector

$$\omega = \frac{2\mu_B B}{\hbar} \dots \text{Larmor frequency}$$

$$\begin{aligned}
 \hat{U} &= e^{i \frac{\omega t}{2} \hat{n} \cdot \vec{\sigma}} = \cos\left(\frac{\omega t}{2} \hat{n} \cdot \vec{\sigma}\right) + i \sin\left(\frac{\omega t}{2} \hat{n} \cdot \vec{\sigma}\right) = \\
 &= \sum_{k=0}^{\infty} \frac{(-1)^k}{(2k)!} \left(\frac{\omega t}{2} \hat{n} \cdot \vec{\sigma}\right)^{2k} + i \sum_{k=0}^{\infty} \frac{(-1)^k}{(2k+1)!} \left(\frac{\omega t}{2} \hat{n} \cdot \vec{\sigma}\right)^{2k+1} = \\
 &= \mathbb{1} \sum_{k=0}^{\infty} \frac{(-1)^k}{(2k)!} \left(\frac{\omega t}{2}\right)^{2k} + i \hat{n} \cdot \vec{\sigma} \sum_{k=0}^{\infty} \frac{(-1)^k}{(2k+1)!} \left(\frac{\omega t}{2}\right)^{2k+1} \\
 &= \mathbb{1} \cdot \cos \frac{\omega t}{2} + i \hat{n} \cdot \vec{\sigma} \sin \frac{\omega t}{2}
 \end{aligned}$$

$$\begin{aligned}
 \vec{n} = n_z: & \quad \hat{n} = \hat{\sigma}_z \Rightarrow \hat{n} \cdot \vec{\sigma} = \hat{\sigma}_z \\
 &= \mathbb{1} \cos \frac{\omega t}{2} + i \hat{\sigma}_z \sin \frac{\omega t}{2} = \begin{pmatrix} \cos \frac{\omega t}{2} + i \sin \frac{\omega t}{2} & 0 \\ 0 & \cos \frac{\omega t}{2} - i \sin \frac{\omega t}{2} \end{pmatrix} = \\
 &= \begin{pmatrix} e^{i \frac{\omega t}{2}} & 0 \\ 0 & e^{-i \frac{\omega t}{2}} \end{pmatrix}
 \end{aligned}$$

Representation on Bloch Sphere:



$$\begin{aligned}
 |\psi(t)\rangle &= U_z |\psi(0)\rangle = \begin{pmatrix} e^{+i\omega t/2} & 0 \\ 0 & e^{-i\omega t/2} \end{pmatrix} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \\
 &= \left(e^{i\omega t/2} |0\rangle\langle 0| + e^{-i\omega t/2} |1\rangle\langle 1| \right) (|0\rangle + |1\rangle) = \\
 &= e^{i\omega t/2} |0\rangle + e^{-i\omega t/2} |1\rangle = e^{i\omega t/2} (|0\rangle + e^{-i\omega t} |1\rangle)
 \end{aligned}$$

Question: In the experiment, how do you implement a phase flip gate?

- a) control the interaction time t
- b) control the strength $|B|$ of the magnetic field
- c) both
- d) none of the above

Question: What is the evolution of the state $|0\rangle$ (on the Bloch sphere), when $\vec{B} = B_z \hat{z}$?

$$\uparrow \vec{B} = \vec{\mu}$$

- a) remains constant
- b) obtains a phase $e^{-i\omega t/2}$
- c) none of the above

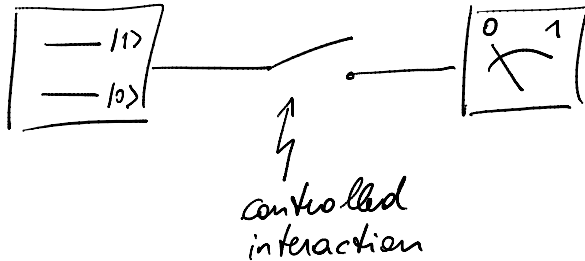
I. 5. Quantum Measurement- Read Out

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generic measurement setup

closed quantum system (QS)

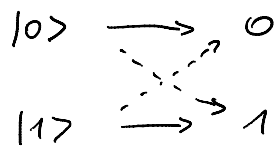
measurement apparatus



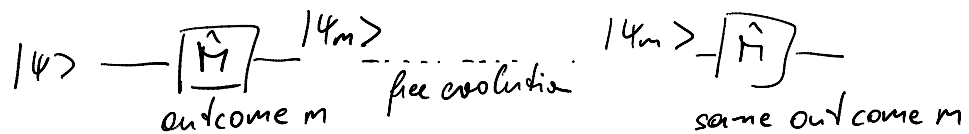
goal: faithful reconstruction of qubit state

What properties should an ideal measurement apparatus have?

- high ON/OFF ratio: no interaction between MA and QS when OFF, strong interaction when ON
- high fidelity of mapping of QS to MA state (no misidentifications)



- fast MA in comparison to decoherence/relaxation
- quantum non-demolition (QND): repeatability of measurement with same outcome



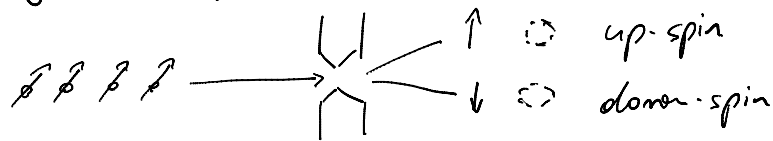
(Projective) Measurement

described by an observable (Hermitian operator) \hat{M} , e.g.

\hat{S}_z ... spin along z-axis

\hat{S}^2 ... modulus square of spin

e.g. Stern-Gerlach apparatus



$$\hat{S}_z = \frac{\hbar}{2} \sigma_z$$

$$\sigma_z = 1 \cdot |1 \times 1\rangle - 1 |0 \times 0\rangle = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

can be decomposed into projection operators \hat{P}_m

$$\hat{M} = \sum_m \hat{P}_m = \sum_m |m\rangle \langle m|$$

↑
eigenvalues m

P_m corresponds to asking the question:

'Is the system in state $|m\rangle$?'

$$P_{\uparrow} |\uparrow\rangle = |\uparrow\rangle \langle \uparrow | \uparrow \rangle = 1 |0\rangle \text{ answer 'yes'}$$

$$P_{\uparrow} |\downarrow\rangle = |\uparrow\rangle \langle \uparrow | \downarrow \rangle = 0 \text{ answer 'no'}$$

Probability, to find outcome m is given by

$$P(m) = \langle \psi | P_m | \psi \rangle$$

State after measurement is given by

$$|\psi_m\rangle = \frac{P_m |\psi\rangle}{\sqrt{\langle \psi | P_m | \psi \rangle}} = \frac{P_m |\psi\rangle}{\sqrt{P_m}} \leftarrow \text{Normalization}$$

e.g. probability to find the value $+1$ when measuring σ_z in superposition state $|\psi\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$

$$P(+1) = \langle \psi | P_1 | \psi \rangle = \langle \psi | 1 \times 1 | \psi \rangle = \\ = \frac{1}{2} (\langle 0 | + \langle 1 |) (1 \times 1) (|0\rangle + |1\rangle) = \frac{1}{2}$$

$$\text{state after measurement is } \frac{P_1 |\psi\rangle}{\sqrt{P_m}} = \frac{1 \times 1 \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)}{\frac{1}{\sqrt{2}}} = |1\rangle \\ \Rightarrow \text{normalized state after measurement}$$

Expectation value: $\langle \hat{M} \rangle = \langle \psi | \hat{M} | \psi \rangle =$

$$= \langle \psi | \sum_m P_m | \psi \rangle =$$

$$= \sum_m \langle \psi | m \times m | \psi \rangle =$$

$$= \sum_m \underbrace{|\langle \psi | m \rangle|^2}_{P(m)}$$

e.g.: σ_z for groundstate

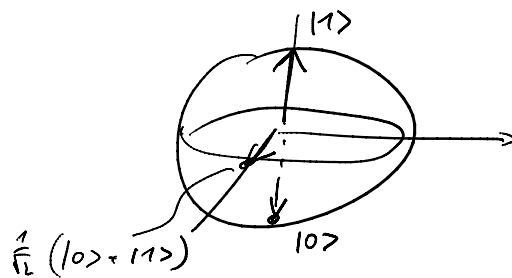
$$\langle 0 | \sigma_z | 0 \rangle = \langle 0 | (1 \times 1 | - |0\rangle \langle 0|) | 0 \rangle \\ = \langle \overset{0}{0} | 1 \rangle \langle 1 | 0 \rangle - \langle \overset{1}{0} | 0 \rangle \langle 0 | 0 \rangle \\ = -1$$

σ_z for excited state $\langle 1 | \sigma_z | 1 \rangle = -1$

σ_z for superpos. state $|\psi\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$

$$\begin{aligned}\langle \psi | \sigma_z | \psi \rangle &= \frac{1}{\sqrt{2}} (\langle 0 | + \langle 1 |) \sigma_z (|0\rangle + |1\rangle) \frac{1}{\sqrt{2}} \\ &= \frac{1}{2} (\langle 0 | + \langle 1 |) (1 \times 1 - 1 \times 0) (|0\rangle + |1\rangle) \\ &= \frac{1}{2} (-1 + 1) = 0\end{aligned}$$

σ_z measures projection of Bloch vector onto z-axis



Questions:

Is the measurement of σ_z sufficient to determine the state completely? (State tomography, see exercises)

Can you perform all required measurements simultaneously using a single particle? ($[\hat{\sigma}_i, \hat{\sigma}_j] \neq 0 \ \forall i \neq j$!)