Quantum computing with trapped ions

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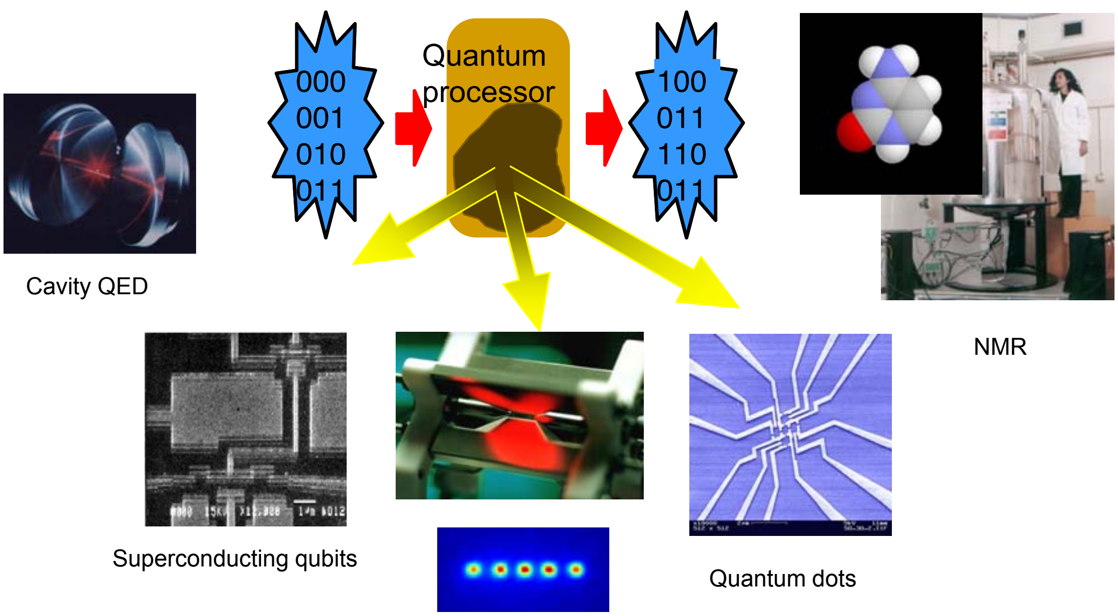
- Basics of ion trap quantum computing
- Measuring a density matrix
- Quantum gates

€ FWF SFB   SCALA QGATES   Industrie Tirol   IQI GmbH
FWF   bm:bwk $

Zürich, Dec 8th 2008

Which technology ?



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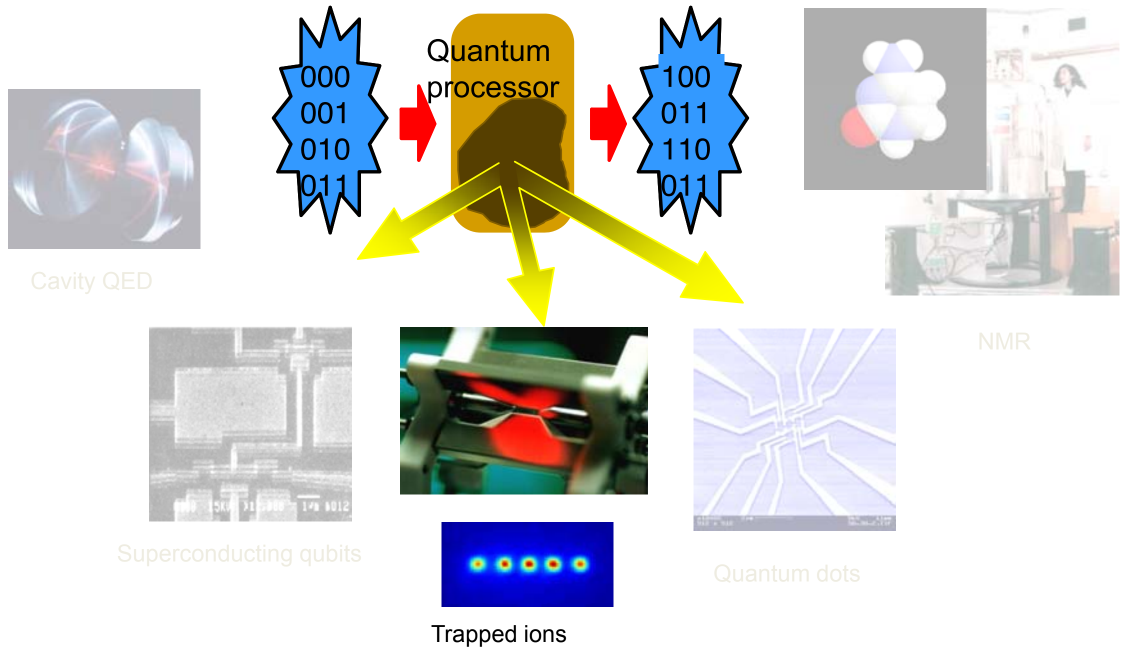


Why trapped ions ?

Good things about ion traps:

- Ions are excellent quantum memories; single qubit coherence times > 10 minutes have been demonstrated

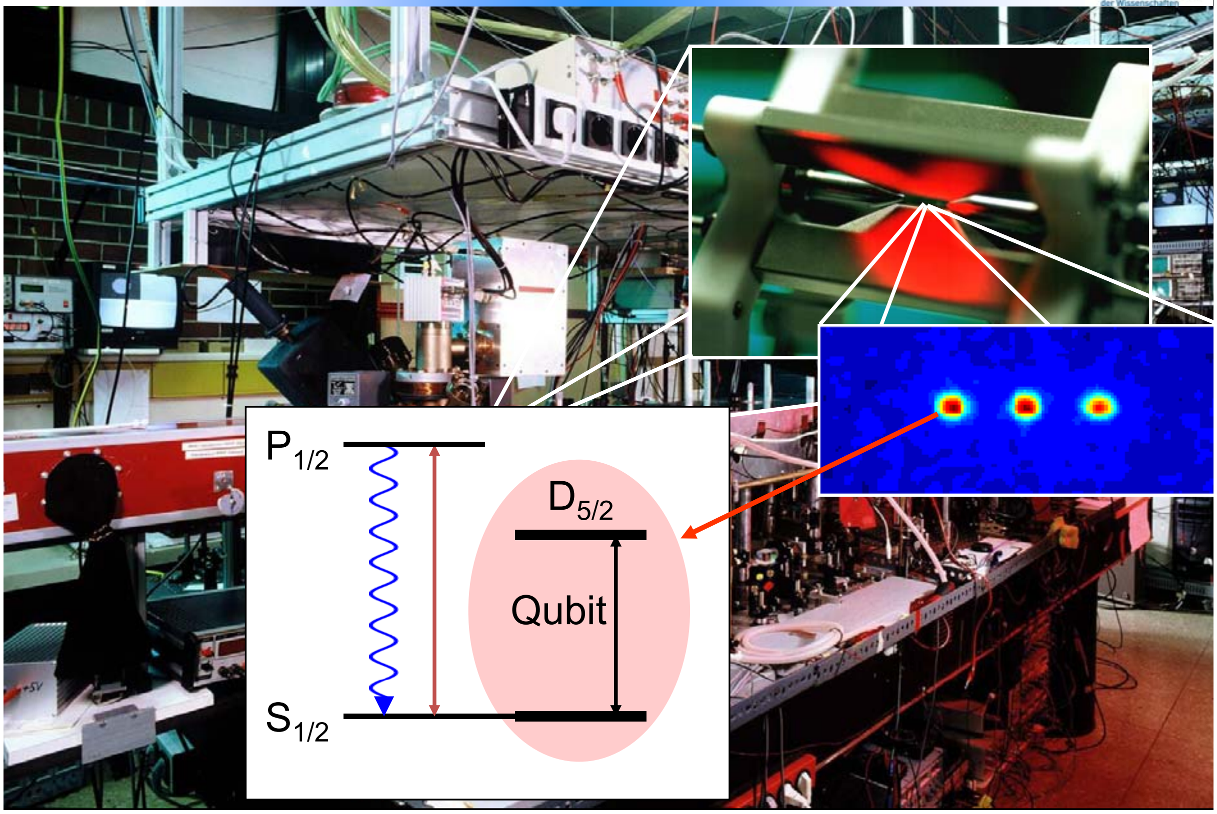
- Ions can be controlled very well

- Many ideas to scale ion traps

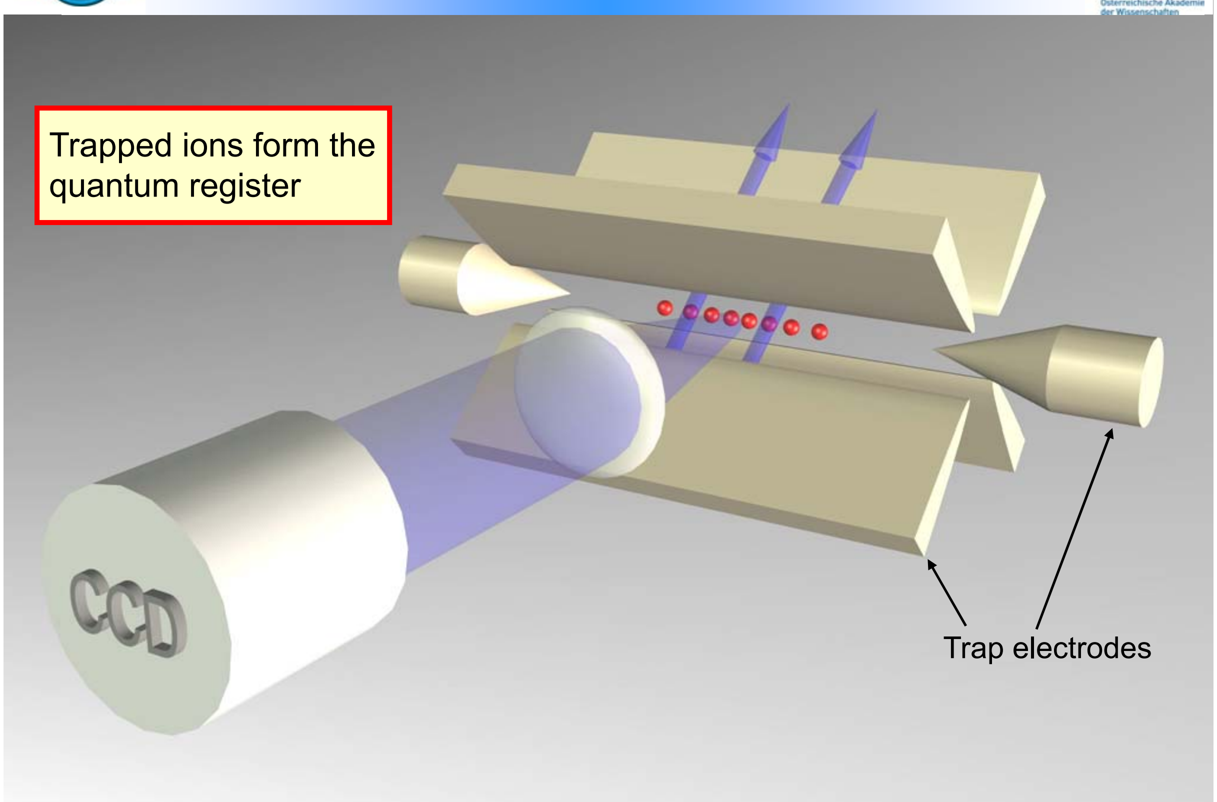
Bad things about ion traps:

- Slow (~1 MHz)
- Technically demanding

The hardware



Innsbruck quantum processor



DiVincenzo criteria

I. Scalable physical system, well characterized qubits

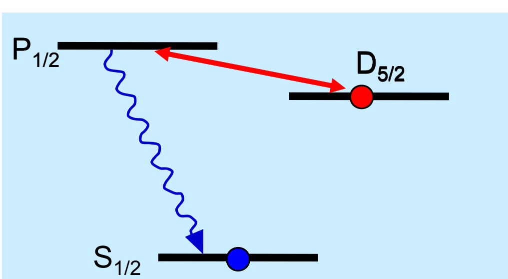
II. Ability to initialize the state of the qubits

III. Long relevant coherence times, much longer than gate operation time

IV. “Universal” set of quantum gates

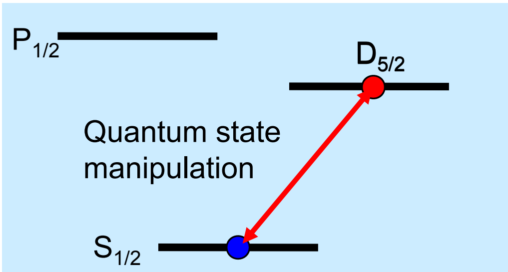
V. Qubit-specific measurement capability

Experimental procedure



1. Initialization in a pure quantum state

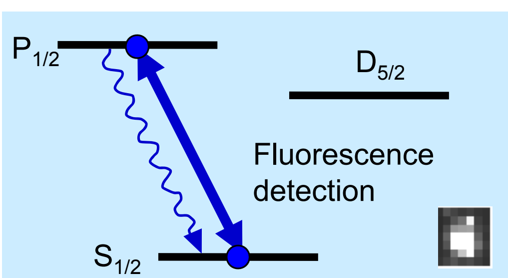
Experimental procedure



1. Initialization in a pure quantum state

2. Quantum state manipulation on $S_{1/2} – D_{5/2}$ transition

Experimental procedure

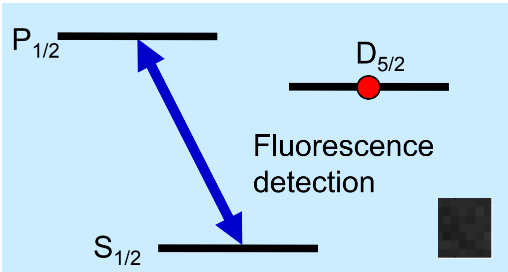


1. Initialization in a pure quantum state:

2. Quantum state manipulation on $S_{1/2} – D_{5/2}$ transition

3. Quantum state measurement by fluorescence detection

Experimental procedure



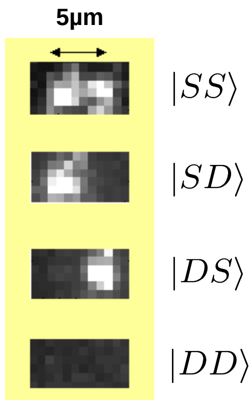
1. Initialization in a pure quantum state:

2. Quantum state manipulation on $S_{1/2} - D_{5/2}$ transition

3. Quantum state measurement by fluorescence detection

Two ions:

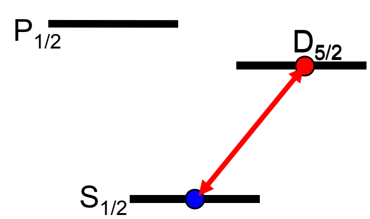
Spatially resolved detection with CCD camera

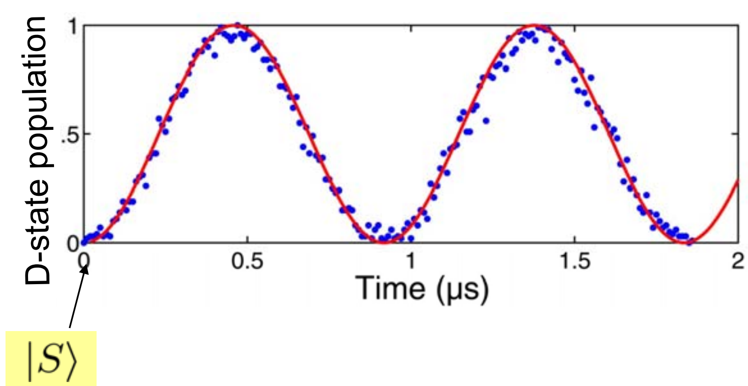


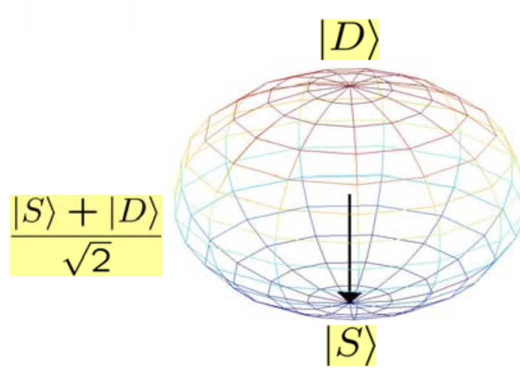
50 experiments / s

Repeat experiments 100-200 times

Rabi oscillations

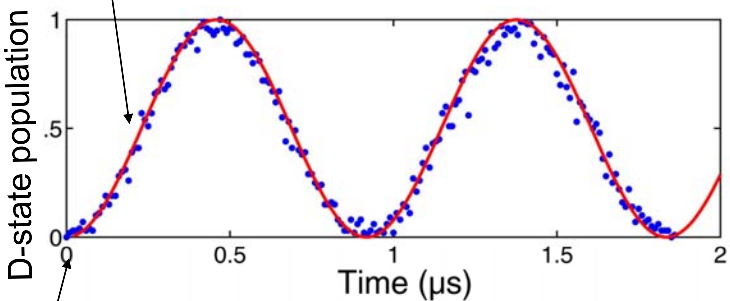




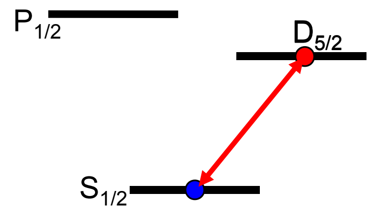


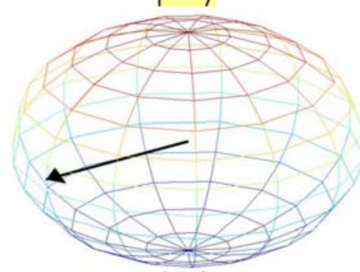
Rabi oscillations

$\frac{|S\rangle + |D\rangle}{\sqrt{2}}$



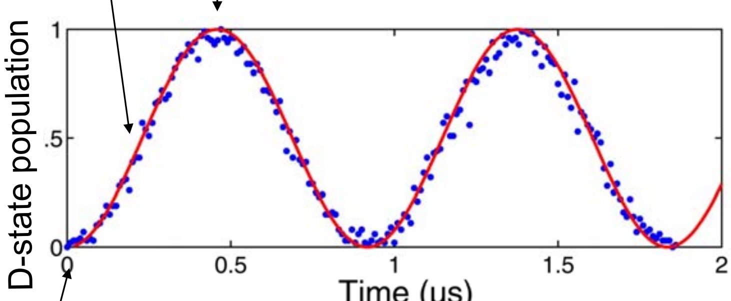
$|S\rangle$



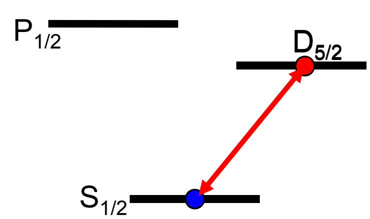


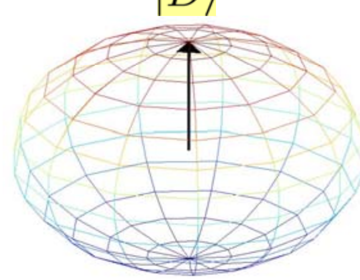
Rabi oscillations

$\frac{|S\rangle + |D\rangle}{\sqrt{2}}$ $|D\rangle$

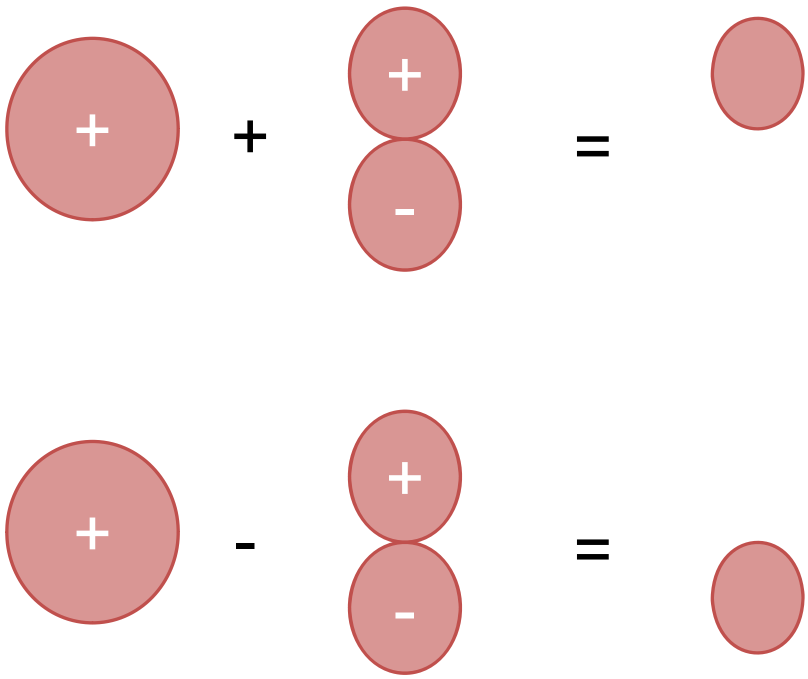


$|S\rangle$





The phase ...



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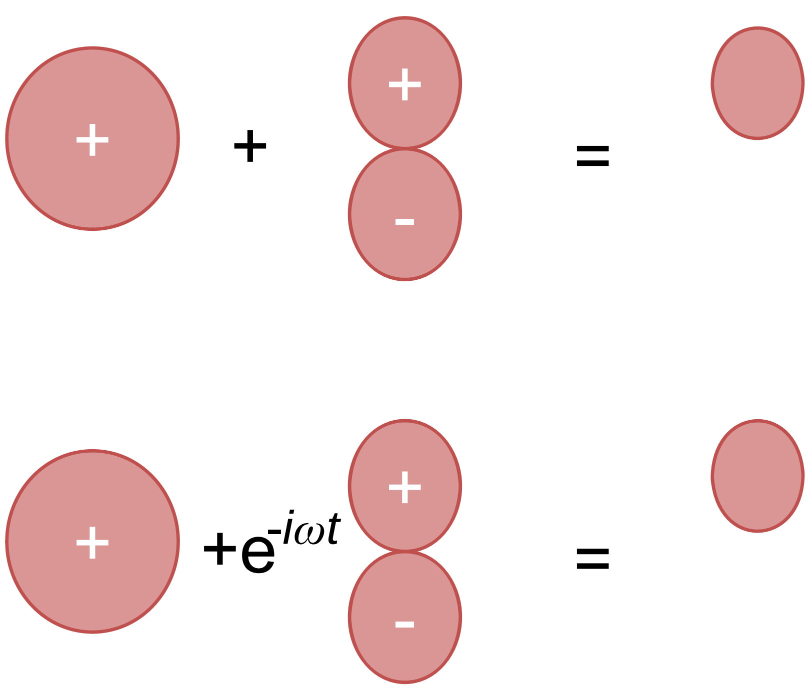

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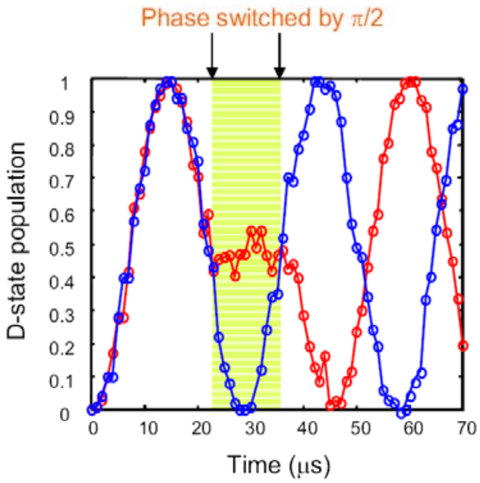

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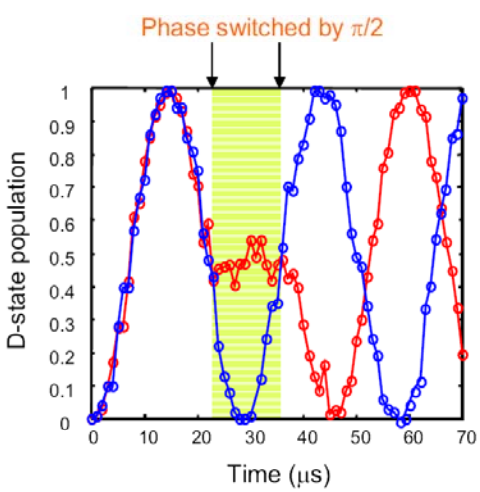
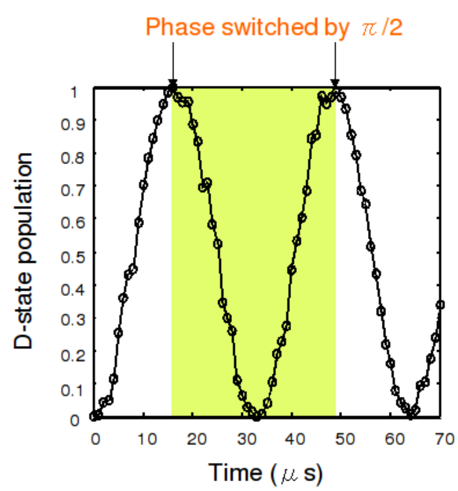
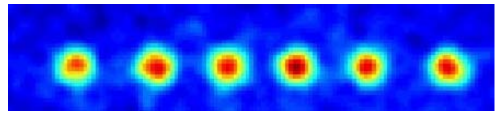
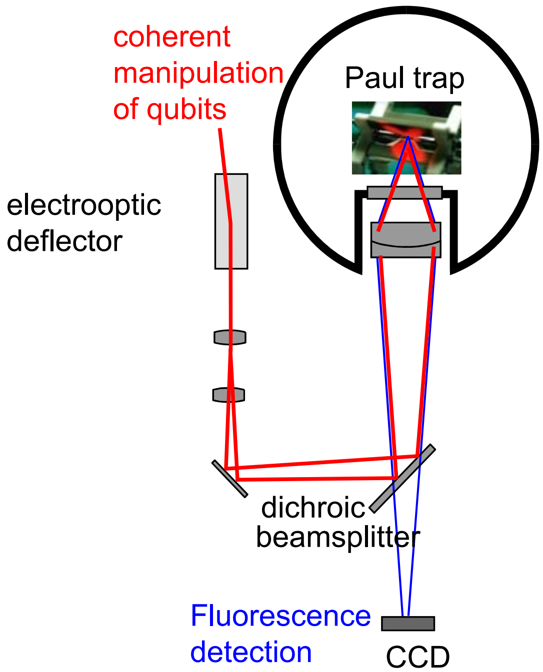
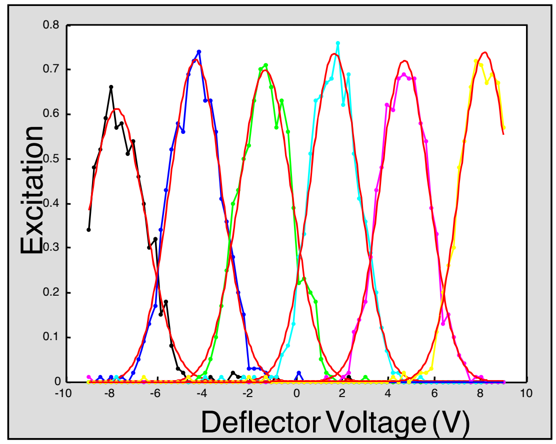




Addressing single qubits







- inter ion distance: ~ 4 μm
- addressing waist: ~ 2 μm

< 0.1% intensity on neighbouring ions

Decoherence mechanisms

Memory errors:

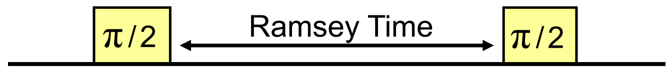
- Bit-flips
- Dephasing

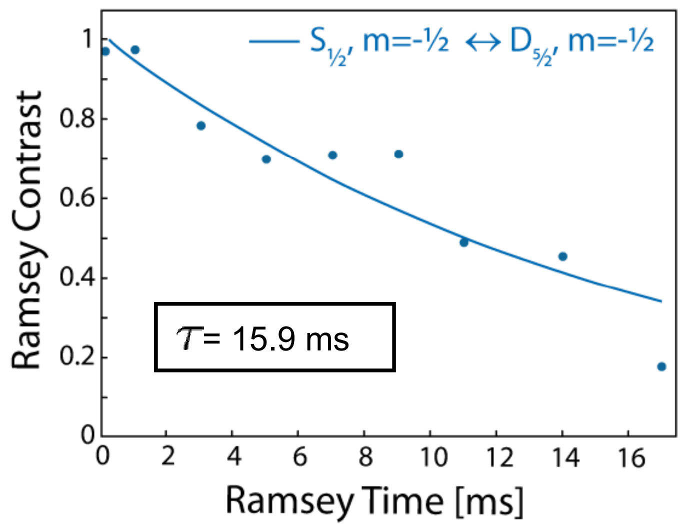
Operationial errors

- technical imperfections …

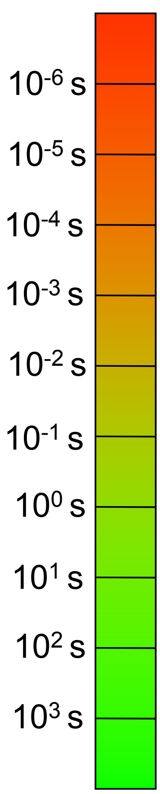
Dephasing of qubits

Ramsey Experiment





Realized time scales



Single qubit gates

Two qubit gates (Geometric phase gates)

Two qubit gates (Cirac-Zoller approach)

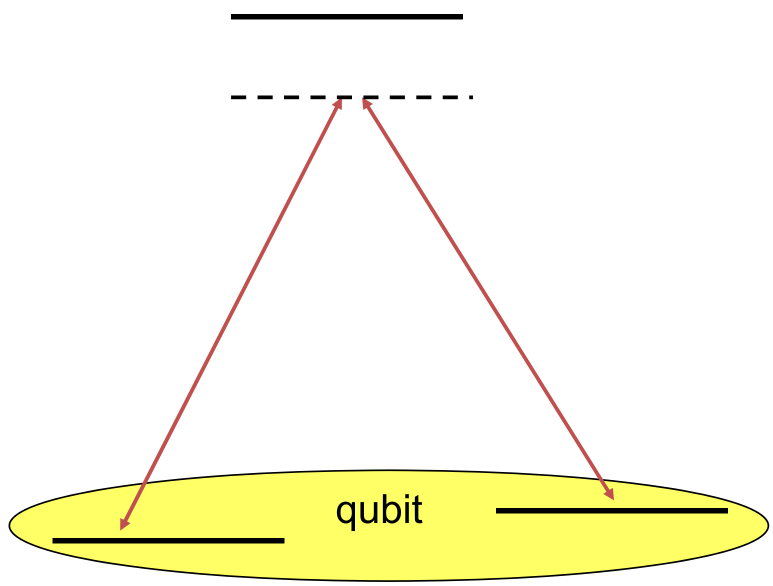
Single qubit coherence (magnetic field sensitive)

Coherence of the motion

Long lived qubits

Raman transitions:

Excited state

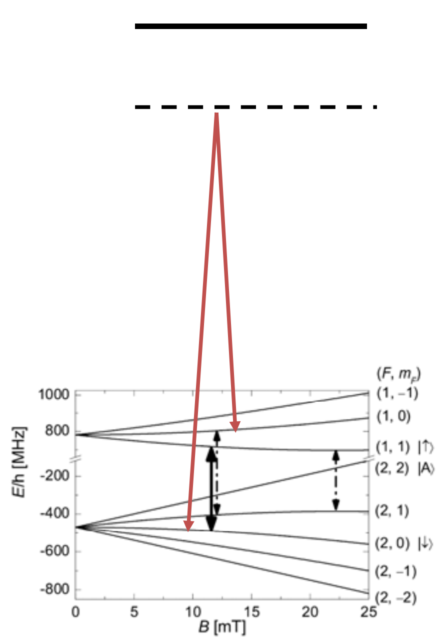


Ground state

Long lived qubits

Raman transitions:

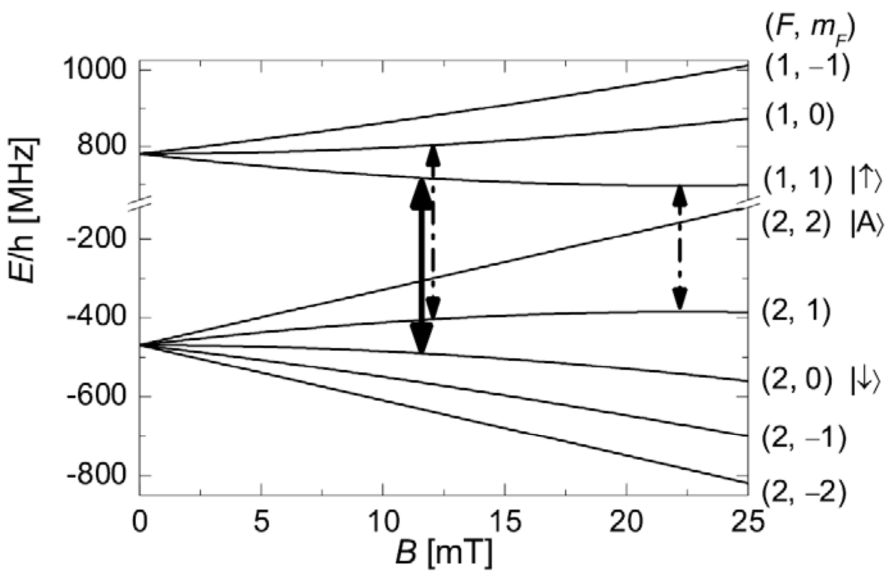
Excited state



Ground state

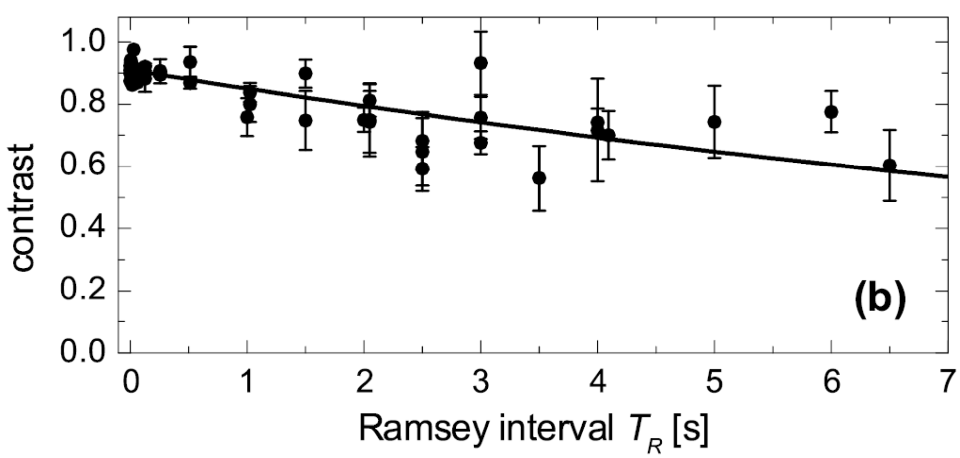
Long lived qubits

Level scheme of $^9Be^+$:



From: C. Langer *et al.*, PRL **95**, 060502 (2005), NIST

Long lived qubits



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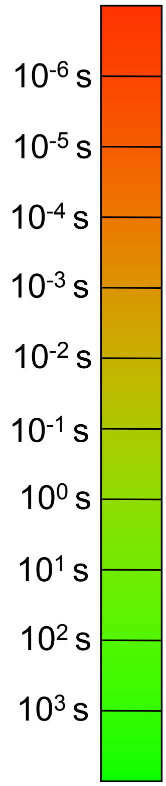


Single qubit gates

Two qubit gates (Geometric phase gates)

Two qubit gates (Cirac-Zoller approach)

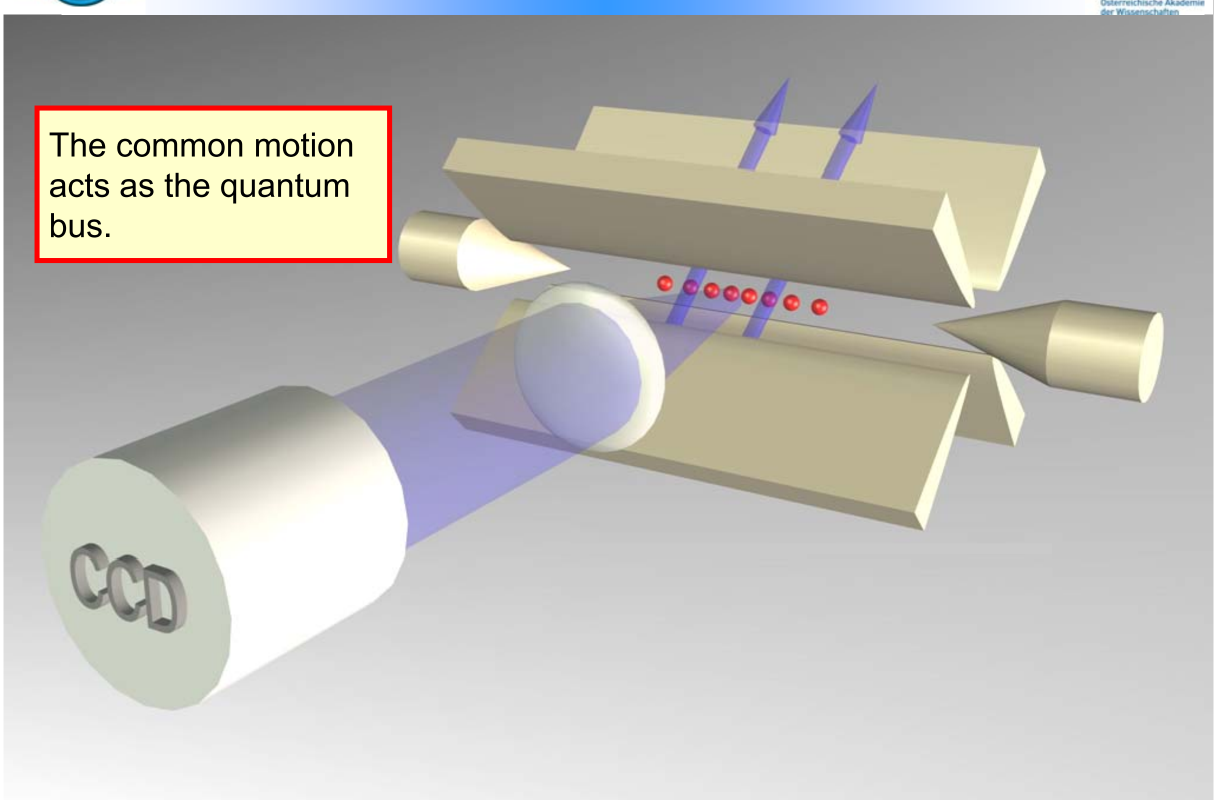
Single qubit coherence (magnetic field sensitive)

Coherence of the motion

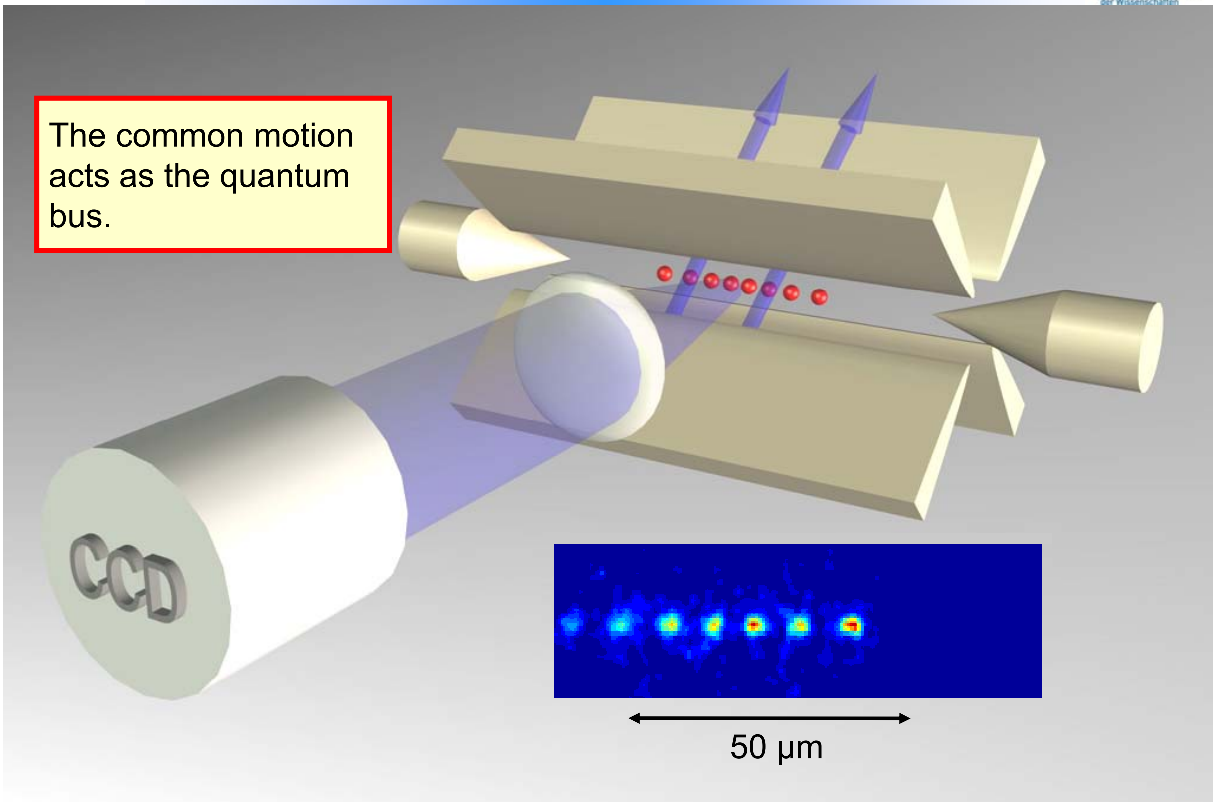
Single qubit coherence (magnetic field insensitive)

Single qubit coherence (magnetic field insensitive + RF drive)

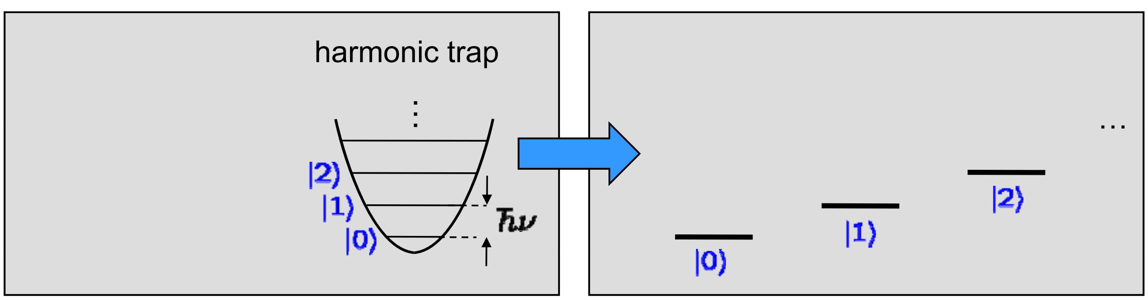
Having the qubits interact



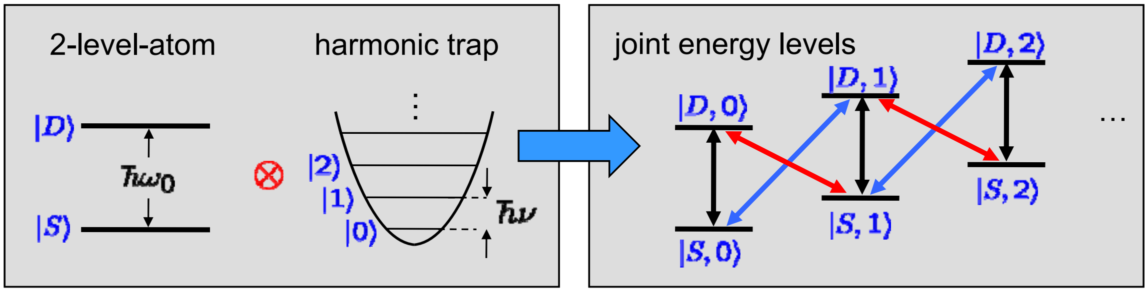
Having the qubits interact

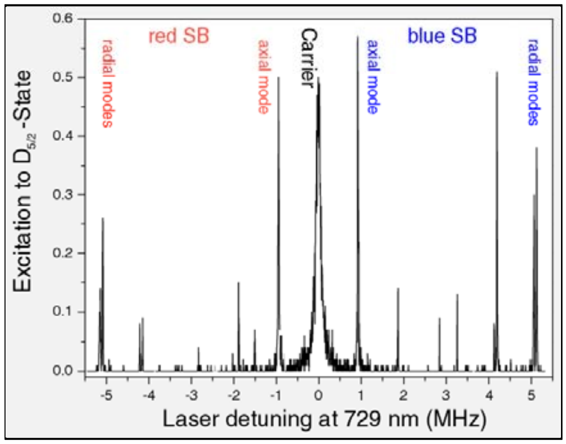


Ion motion

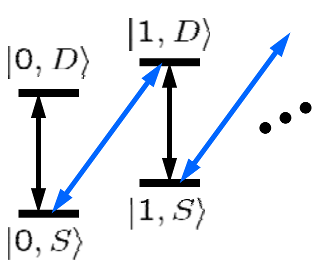


Ion motion

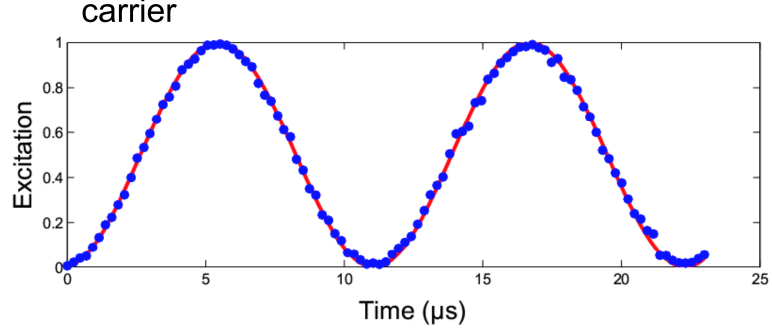




Coherent manipulation

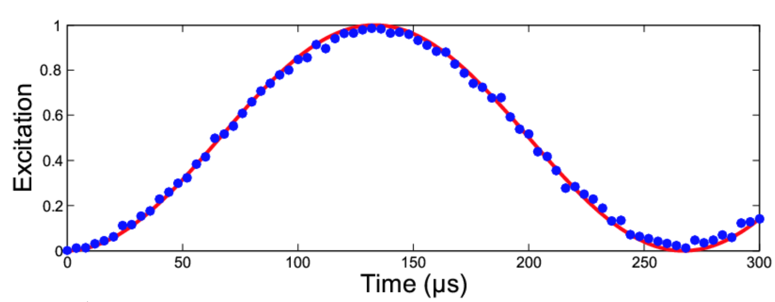


carrier



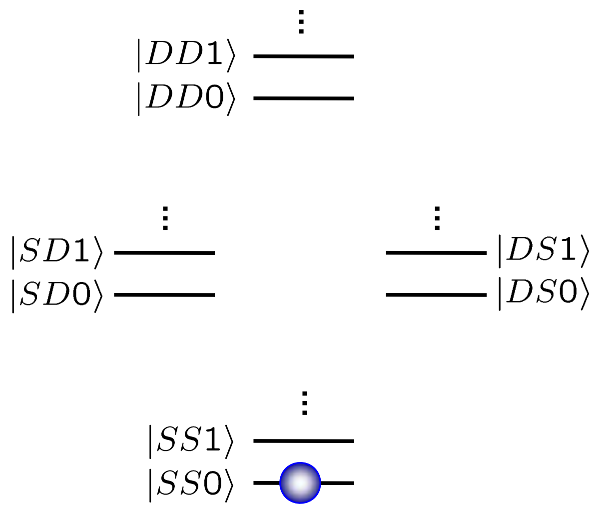
carrier and sideband
Rabi oscillations
with Rabi frequencies

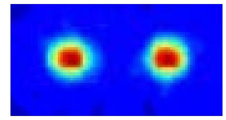
$$\Omega, \quad \eta\Omega$$



$\eta = kx_0$ Lamb-Dicke parameter

Generation of Bell states





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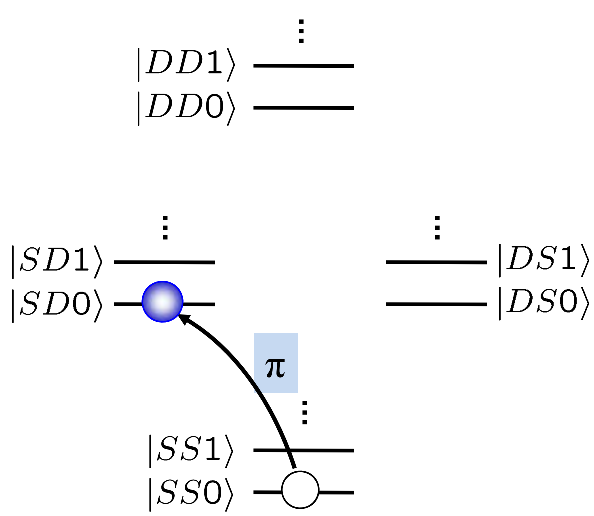

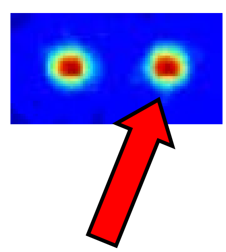

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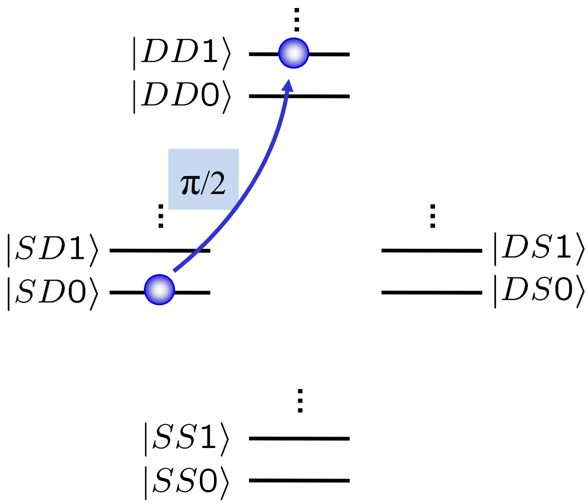

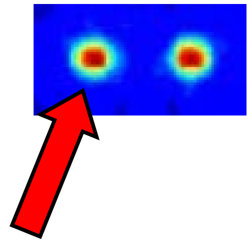

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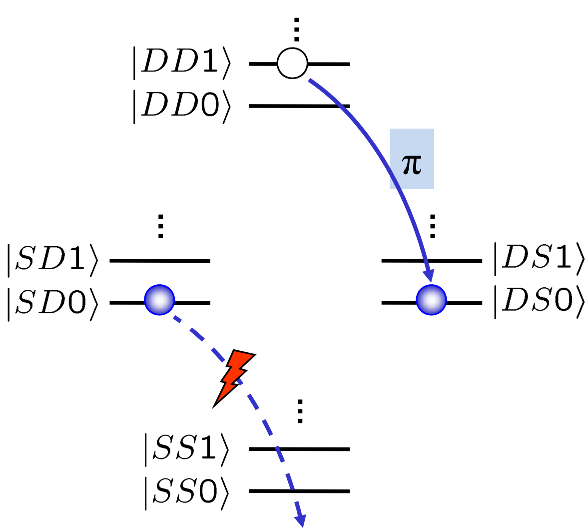

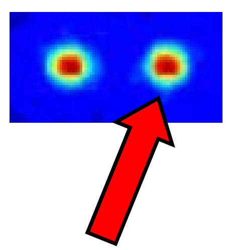

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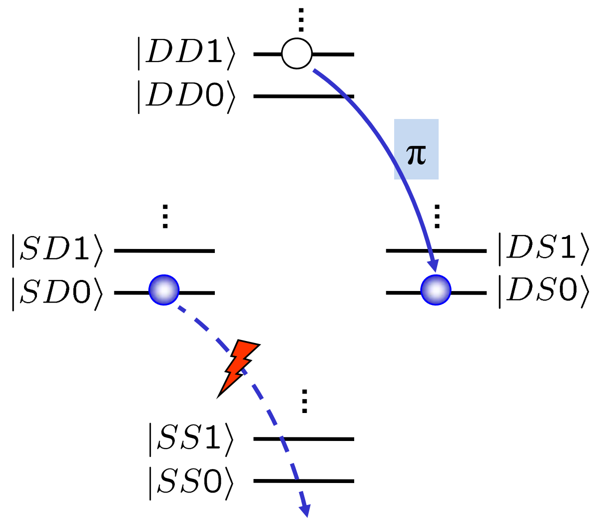
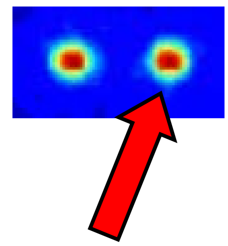




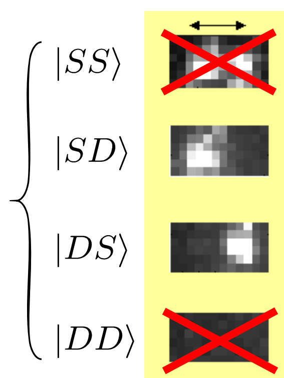
<u>Bell states with atoms</u>

- $^{9}Be^{+}$: NIST (fidelity: 97 %)
- $^{40}Ca^{+}$: Oxford (83%)
- $^{111}Cd^{+}$: Ann Arbor (79%)
- $^{25}Mg^{+}$: Munich
- $^{40}Ca^{+}$: Innsbruck (99%)

Analysis of Bell states

$$|SD\rangle + |DS\rangle$$

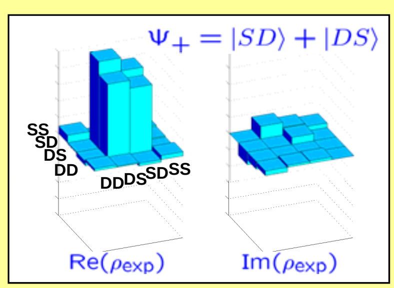
Fluorescence detection with CCD camera:



Coherent superposition or incoherent mixture ?

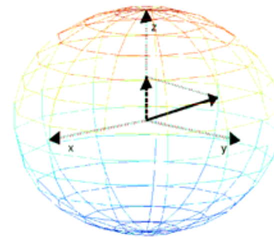
What is the relative phase of the superposition ?

➡ Measurement of the density matrix:



Measuring a density matrix

$|D\rangle$



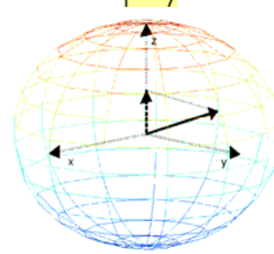
$|S\rangle$

A measurement yields the *z*-component of the Bloch vector
=> Diagonal of the density matrix

$$\rho = \begin{pmatrix} P_S & C - iD \\ C + iD & P_D \end{pmatrix}$$

Measuring a density matrix

$|D\rangle$



$|S\rangle$

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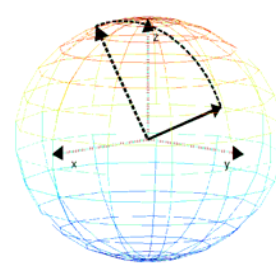
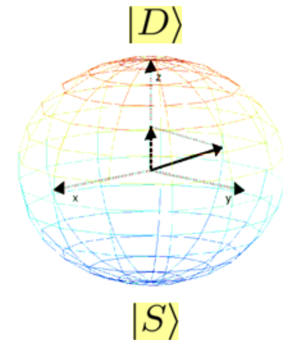
Rotation around the *x*- or the *y*-axis prior to the measurement yields the phase information of the qubit.



Measuring a density matrix



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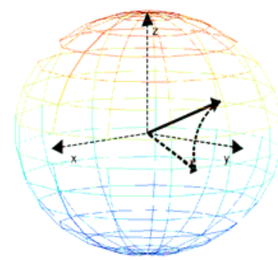
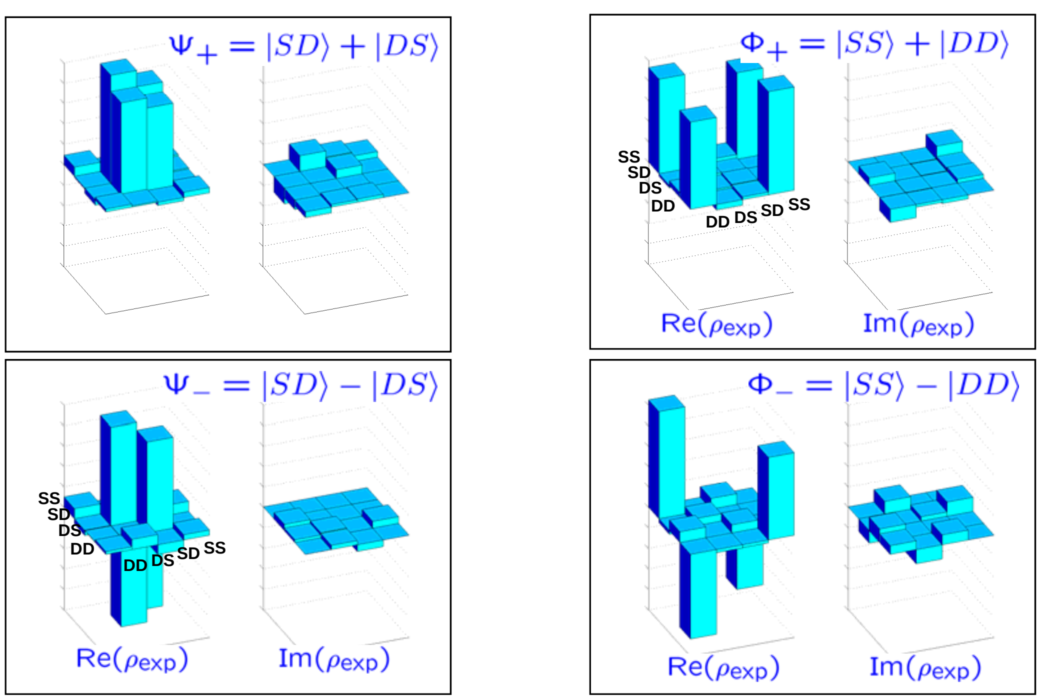
=> coherences of the density matrix

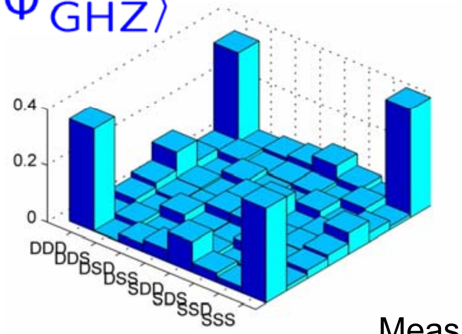


Decoherence properties of qubits

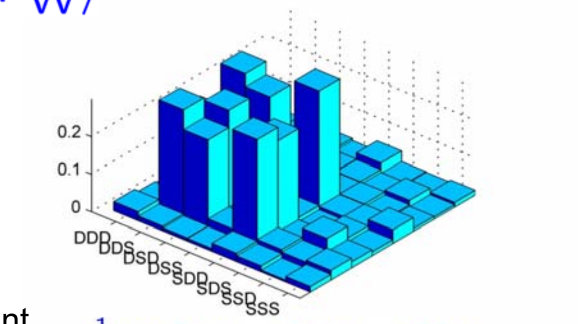


A “real” thought experiment

$|\Psi_{GHZ}\rangle$



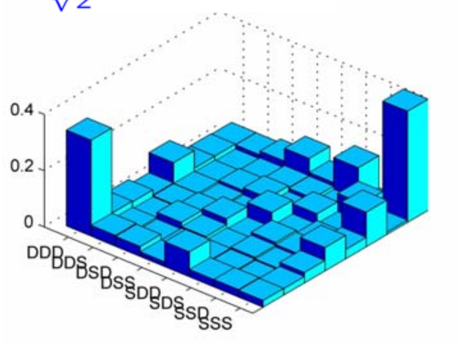
$|\Psi_{W}\rangle$

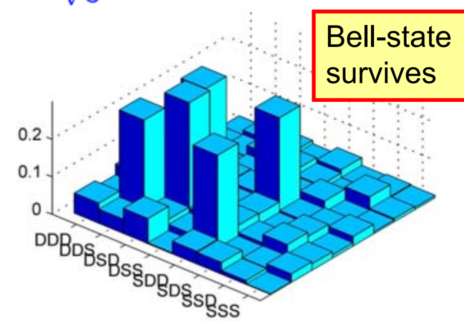


$\frac{1}{\sqrt{2}}(|SSS\rangle + |DDD\rangle)$

Measurement of the center ion

$\frac{1}{\sqrt{3}}(|SDD\rangle + |DSD\rangle + |DDS\rangle)$



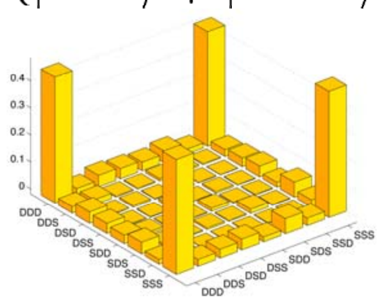


Bell-state survives

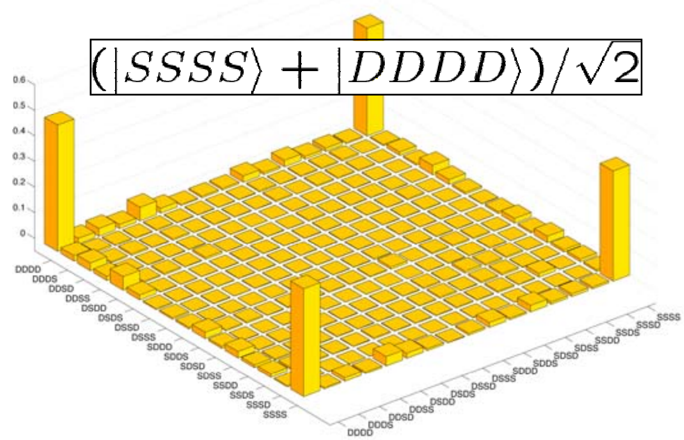
Roos et al., *Science* **304**, 1478 (20

Generalized Bell states

$(|SSS\rangle + |DDD\rangle)/\sqrt{2}$

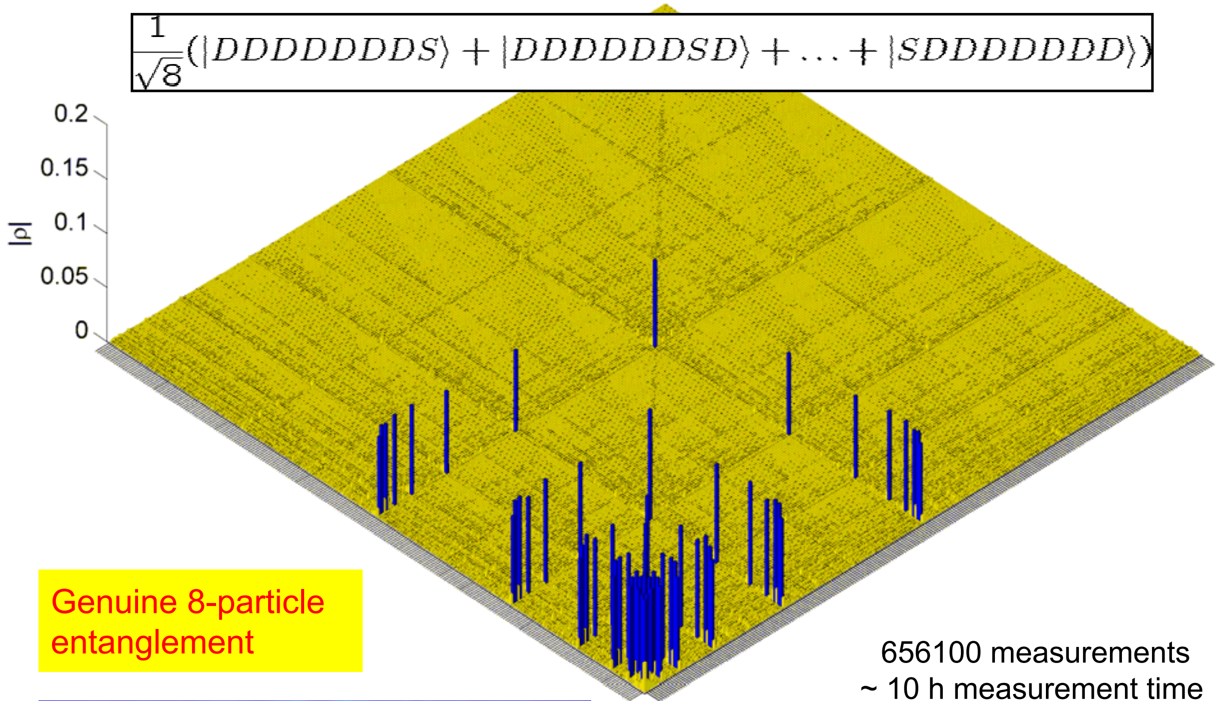


$\boxed{(|SSSS\rangle + |DDDD\rangle)/\sqrt{2}}$

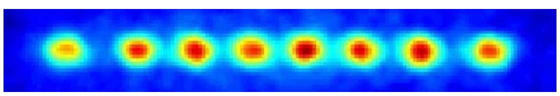


Generalized Bell states

$$\frac{1}{\sqrt{8}}(|DDDDDDDS\rangle + |DDDDDDSD\rangle + \ldots + |SDDDDDDD\rangle)$$



Genuine 8-particle entanglement



656100 measurements
~ 10 h measurement time

Häffner et al., Nature **438**, 643 (2005)

DiVincenzo criteria

I. Scalable physical system, well characterized qubits

II. Ability to initialize the state of the qubits

III. Long relevant coherence times, much longer than gate operation time

IV. “Universal” set of quantum gates

V. Qubit-specific measurement capability

Quantum gates …

Having the qubits interact

VOLUME 74, NUMBER 20 PHYSICAL REVIEW LETTERS 15 MAY 1995

Quantum Computations with Cold Trapped Ions

J. I. Cirac and P. Zoller*

Institut für Theoretische Physik, Universiät Innsbruck, Technikerstrasse 25, A-6020 Innsbruck, Austria

(Received 30 November 1994)

A quantum computer can be implemented with cold ions confined in a linear trap and interacting with laser beams. Quantum gates involving any pair, triplet, or subset of ions can be realized by coupling the ions through the collective quantized motion. In this system decoherence is negligible, and the measurement (readout of the quantum register) can be carried out with a high efficiency.

PACS numbers: 89.80.+h, 03.65.Bz, 12.20.Fv, 32.80.Pj

...allows the realization of a ***universal*** quantum computer !

$$|D\rangle|D\rangle \rightarrow |D\rangle|D\rangle$$

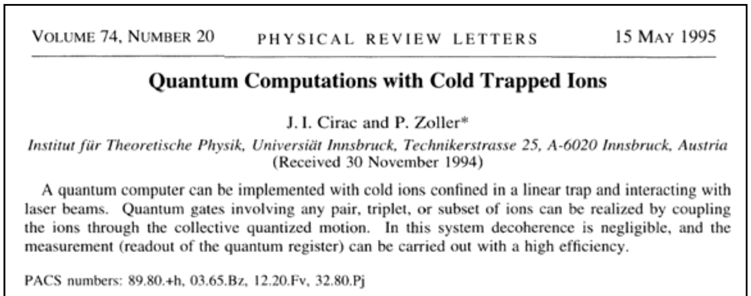
$$|D\rangle|S\rangle \rightarrow |D\rangle|S\rangle$$

$$|S\rangle|D\rangle \rightarrow |D\rangle|S\rangle$$

$$|S\rangle|S\rangle \rightarrow |S\rangle|D\rangle$$

control target

Having the qubits interact



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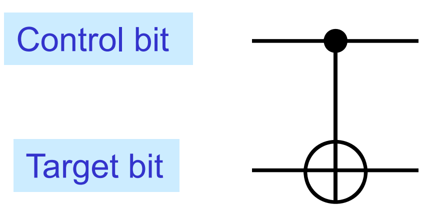
$$|D\rangle|D\rangle \rightarrow |D\rangle|D\rangle$$
$$|D\rangle|S\rangle \rightarrow |D\rangle|S\rangle$$
$$|S\rangle|D\rangle \rightarrow |D\rangle|S\rangle$$
$$|S\rangle|S\rangle \rightarrow |S\rangle|D\rangle$$

control target

Most popular gates:
- Cirac-Zoller gate (Schmidt-Kaler et al., Nature **422,** 408 (2003)).
- Geometric phase gate (Leibfried et al., Nature **422,** 412 (2003)).
- Mølmer-Sørensen gate (Sackett et al., Nature **404**, 256 (2000)).

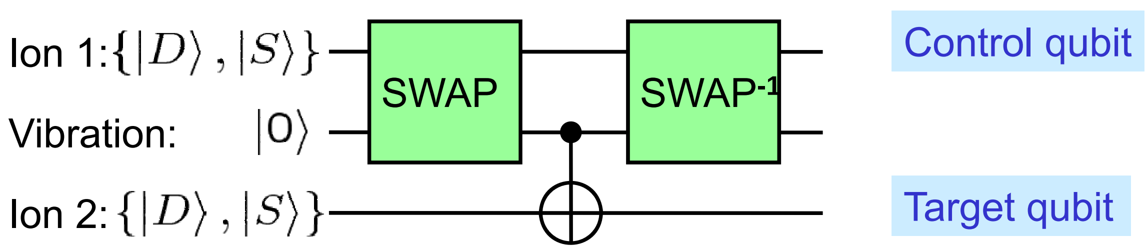
A controlled-NOT operation



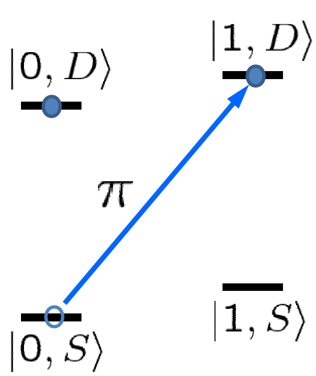
$$|0\rangle|0\rangle \rightarrow |0\rangle|0\rangle$$
$$|0\rangle|1\rangle \rightarrow |0\rangle|1\rangle$$
$$|1\rangle|0\rangle \rightarrow |1\rangle|1\rangle$$
$$|1\rangle|1\rangle \rightarrow |1\rangle|0\rangle$$

Target

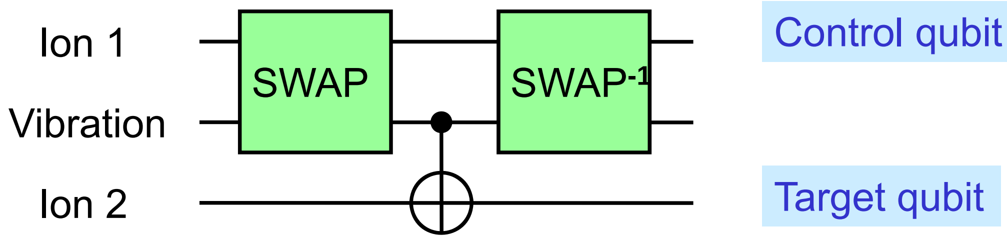
A controlled-NOT operation



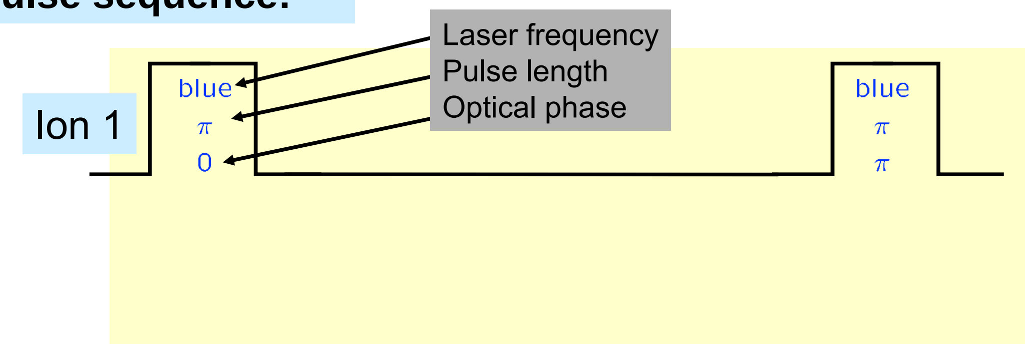
Mapping the qubit to the bus



A controlled-NOT operation

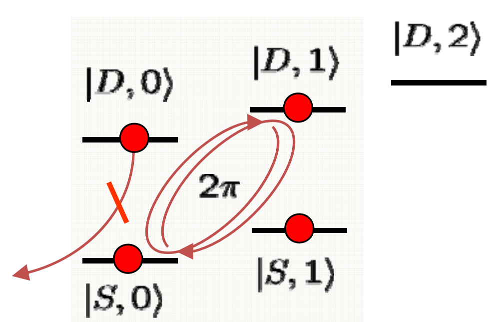


Pulse sequence:

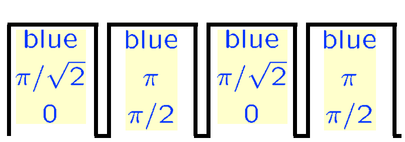


A phase gate

$$
\begin{array}{ccccc}
 & |D,0\rangle & |S,0\rangle & |D,1\rangle & |S,1\rangle \\
U_\Phi = & \multicolumn{4}{c}{\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}}
\end{array}
$$

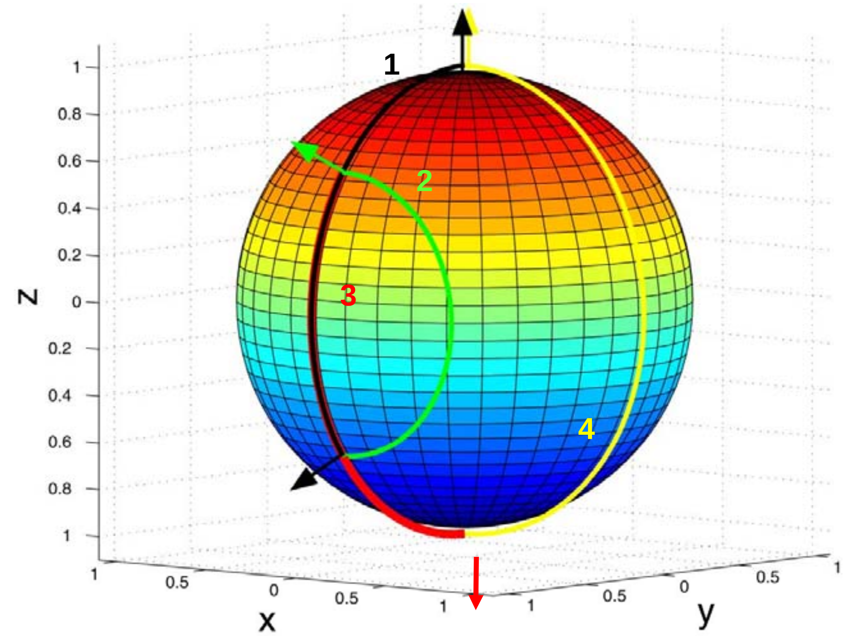


Composite 2π-rotation:

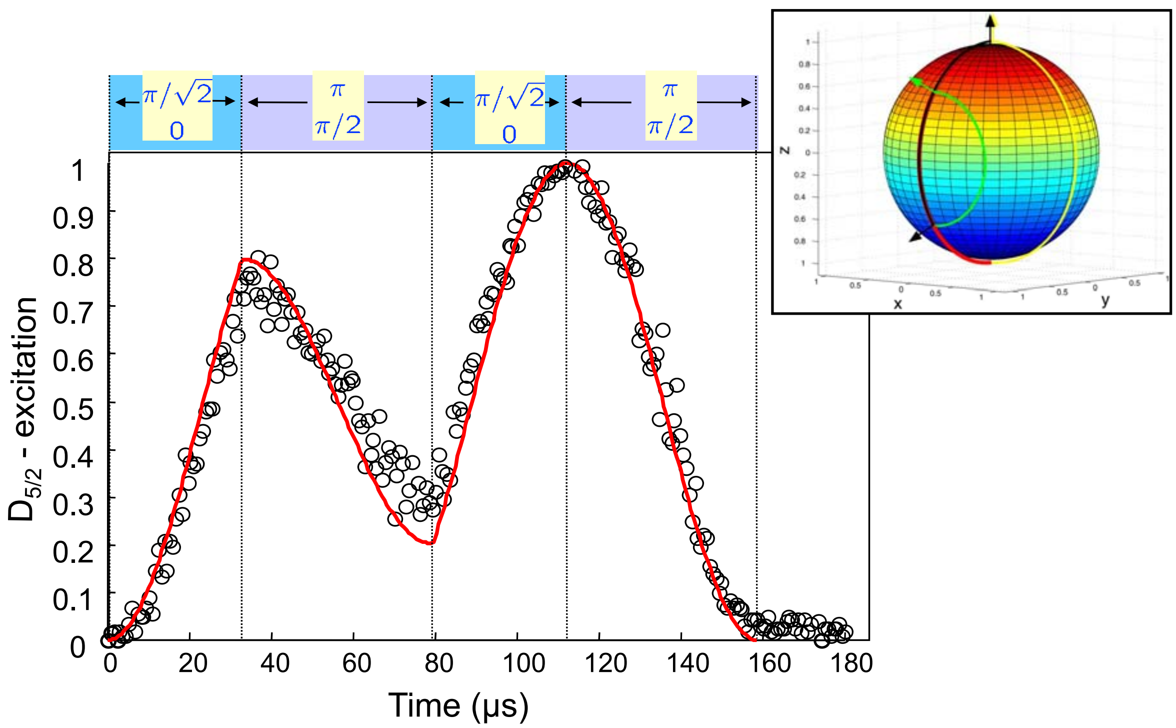


Composite phase gate

$$R^{+}(\pi,\pi/2)\,R^{+}\left(\pi/\sqrt{2},0\right)\,R^{+}(\pi,\pi/2)\,R^{+}\left(\pi/\sqrt{2},0\right)$$

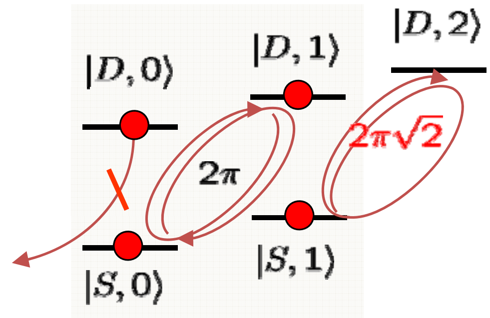


Composite phase gate



A phase gate

$$
\begin{array}{c} \\ U_\Phi = \end{array}
\begin{array}{c}
\begin{array}{cccc} |D,0\rangle & |S,0\rangle & |D,1\rangle & |S,1\rangle \end{array} \\
\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & \boxed{?} \end{pmatrix}
\end{array}
$$



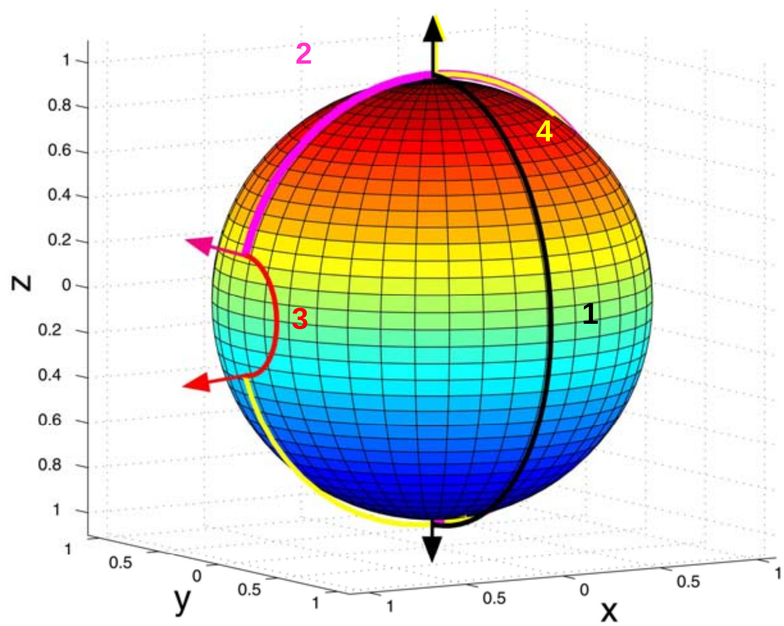
Composite 2π-rotation:

$$a^\dagger |n\rangle = \sqrt{n+1}\,|n+1\rangle$$

blue	blue	blue	blue
$\pi/\sqrt{2}$	π	$\pi/\sqrt{2}$	π
0	$\pi/2$	0	$\pi/2$

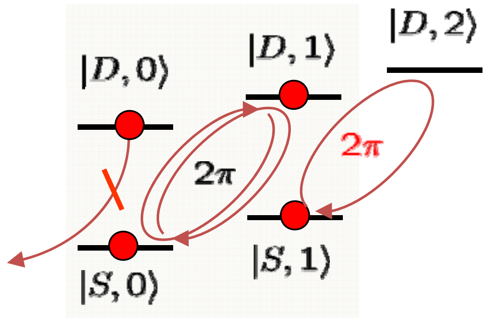
Action on the |1,S> ↔ |2,D> manifold

$$R^+(\pi,\pi/2)\, R^+\left(\pi/\sqrt{2},0\right)\, R^+(\pi,\pi/2)\, R^+\left(\pi/\sqrt{2},0\right)$$



A phase gate

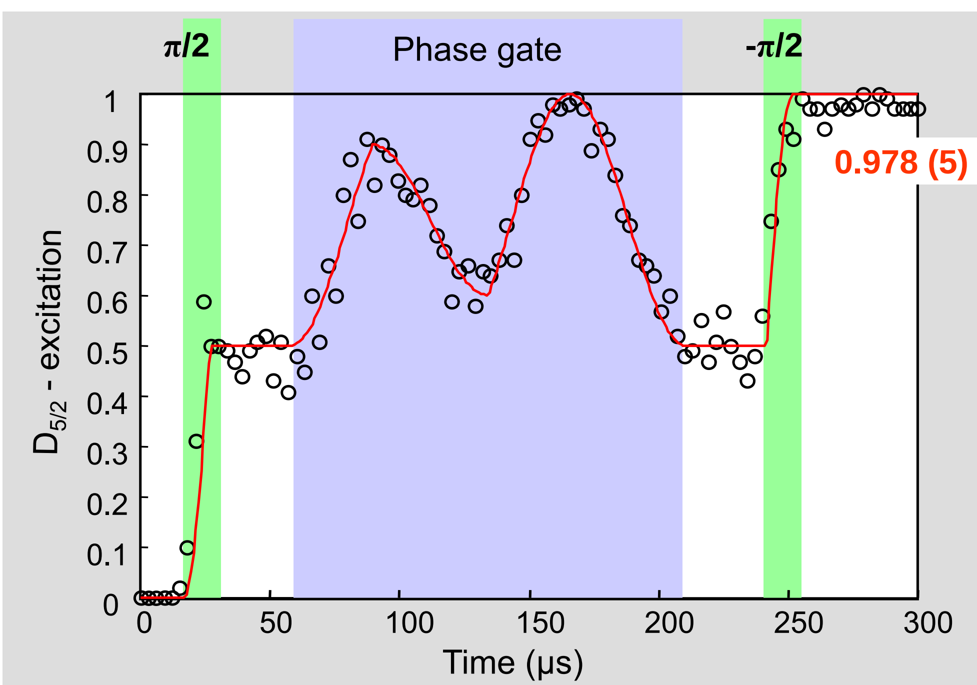
$$
\begin{array}{ccccc}
 & |D,0\rangle & |S,0\rangle & |D,1\rangle & |S,1\rangle \\
U_\Phi = & \multicolumn{4}{c}{\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}}
\end{array}
$$



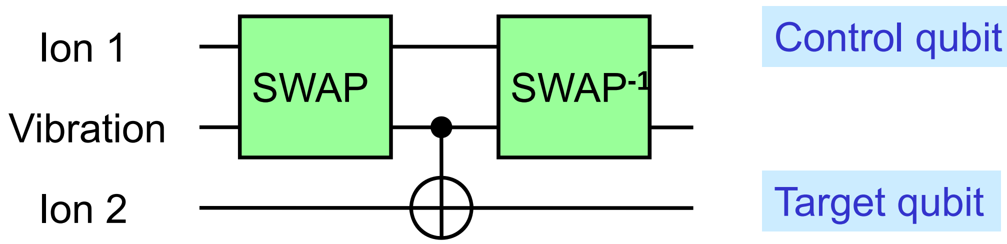
Composite 2π-rotation:

blue	blue	blue	blue
$\pi/\sqrt{2}$	π	$\pi/\sqrt{2}$	π
0	$\pi/2$	0	$\pi/2$

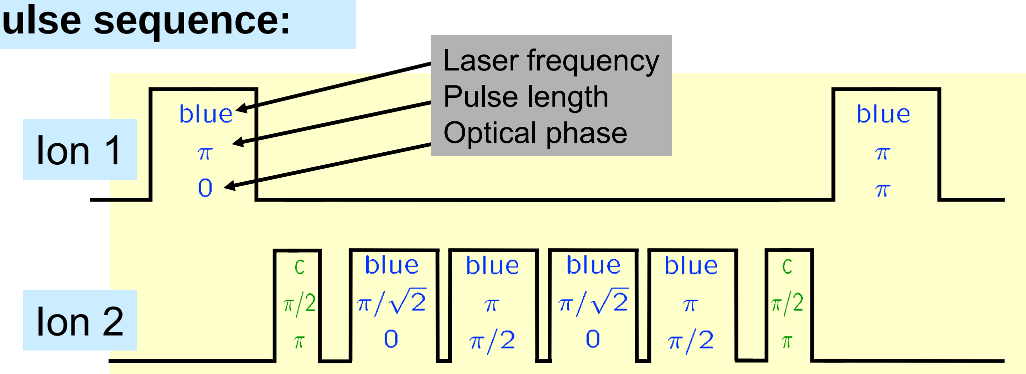
Single-ion CNOT



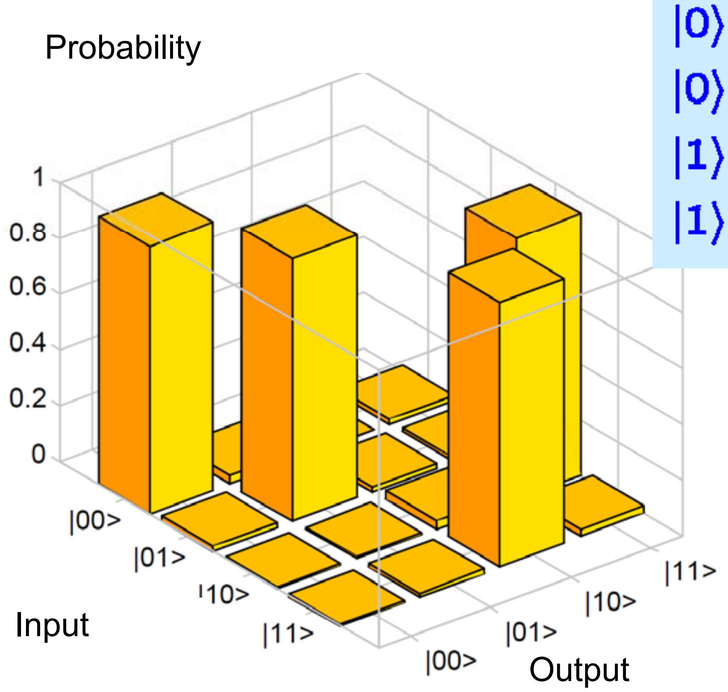
A controlled-NOT operation



Pulse sequence:



Truth table of the CNOT

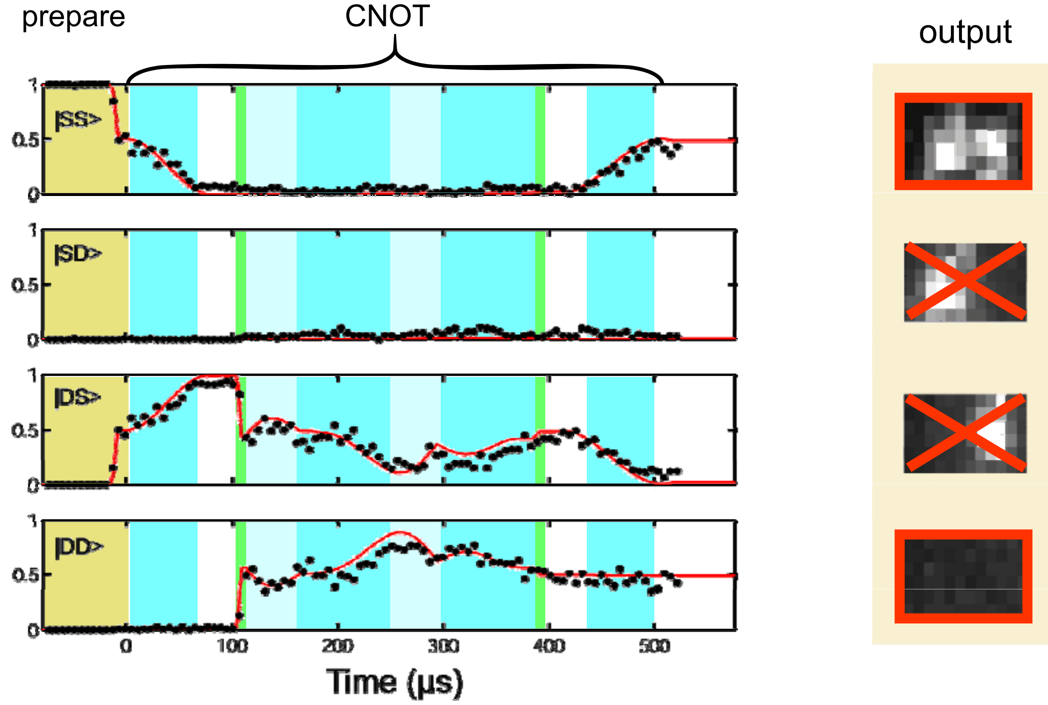


$$|0\rangle|0\rangle \rightarrow |0\rangle|0\rangle$$
$$|0\rangle|1\rangle \rightarrow |0\rangle|1\rangle$$
$$|1\rangle|0\rangle \rightarrow |1\rangle|1\rangle$$
$$|1\rangle|1\rangle \rightarrow |1\rangle|0\rangle$$

Using a CNOT to create a Bell state

$$|SS\rangle \rightarrow |S+D\rangle|S\rangle \quad \longrightarrow \quad |SS\rangle + |DD\rangle$$

prepare CNOT

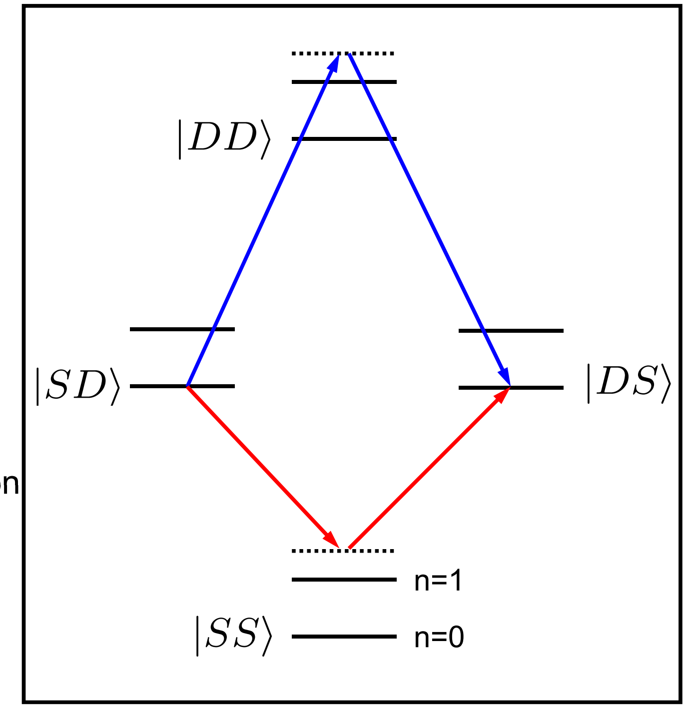


Cirac-Zoller gate

Draw backs of the Cirac-Zoller gate:

- slow (200 trap periods)
- single ion addressing required

Mølmer-Sørensen gate

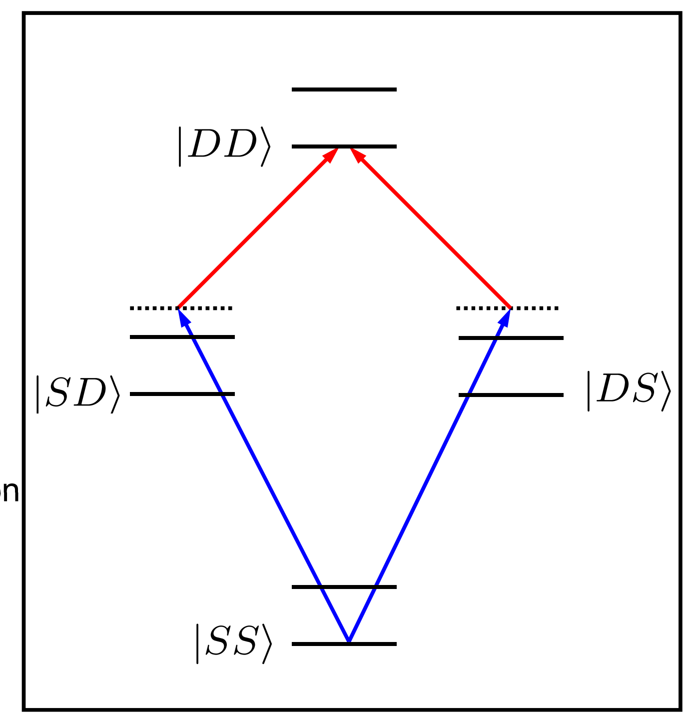


Raman transitions between

$$|SD\rangle \quad \Leftrightarrow \quad |DS\rangle$$

Interaction of two ions via common motion.

Mølmer-Sørensen gate

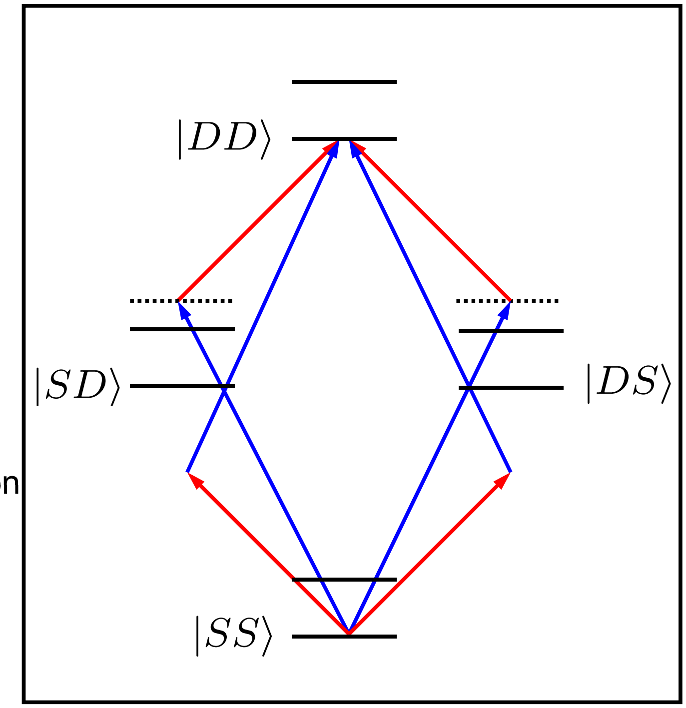


Raman transitions between

$$|SS\rangle \quad \Leftrightarrow \quad |DD\rangle$$

Interaction of two ions via common motion.

Mølmer-Sørensen gate



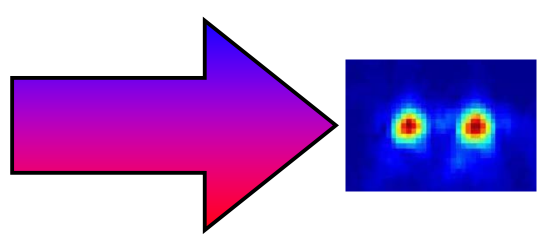
Raman transitions between

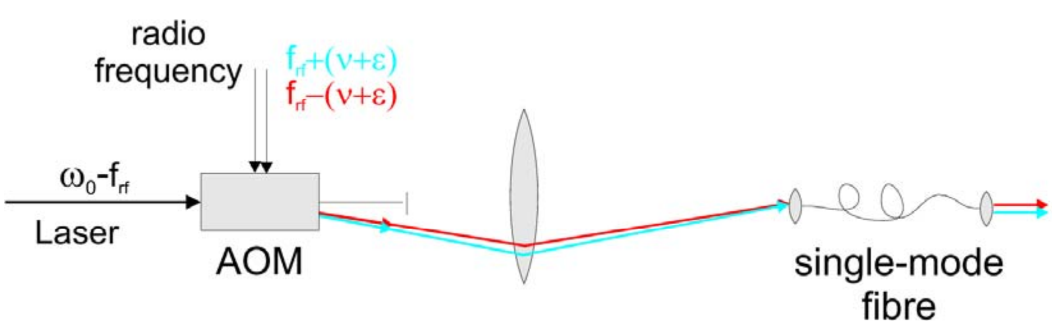
$$|SS\rangle \quad \Leftrightarrow \quad |DD\rangle$$

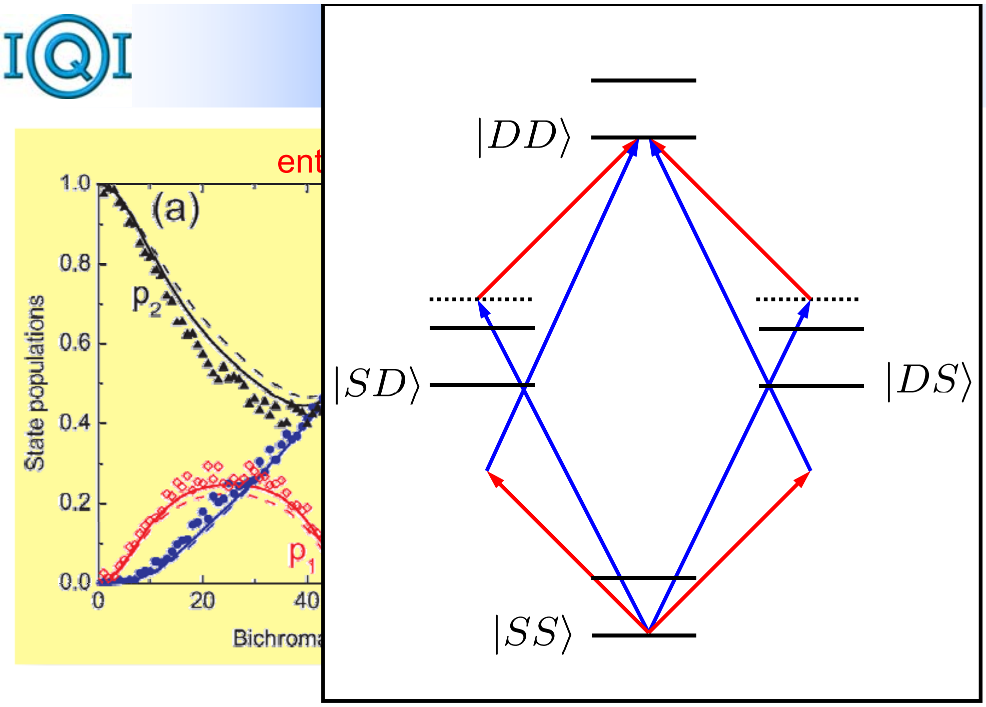
Interaction of two ions via common motion.

Technical realization

bicromatic beam
applied to both ions



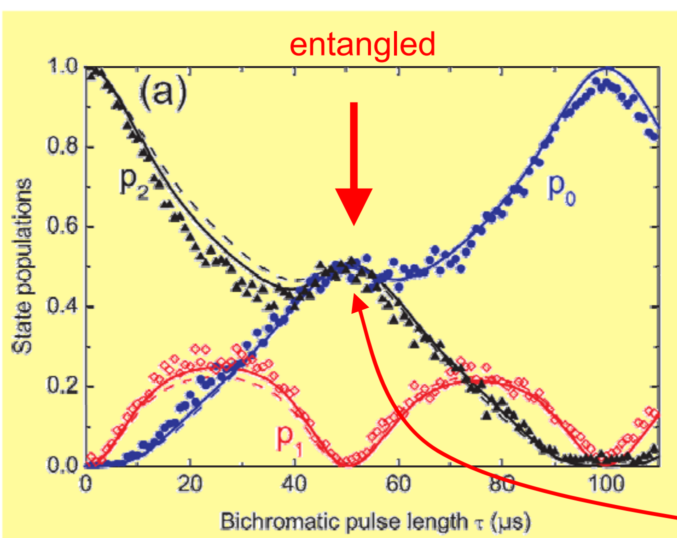




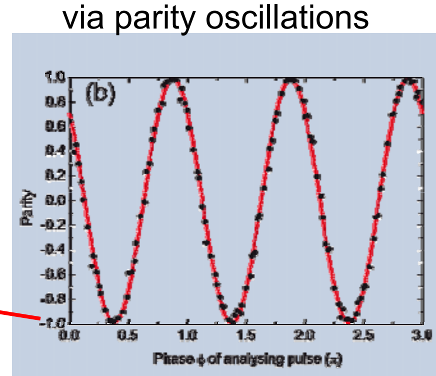
J. Benhelm et al., Nature Physics **4**, 463 (2008)

Theory: C. Roos, NJP **10,** 013002 (2008)

Entangling ions



measure entanglement via parity oscillations



gate duration $51\mu s$

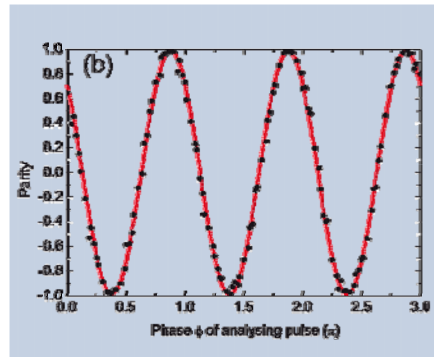
average fidelity: 99.3 (2) %

J. Benhelm et al., Nature Physics **4**, 463 (2008)

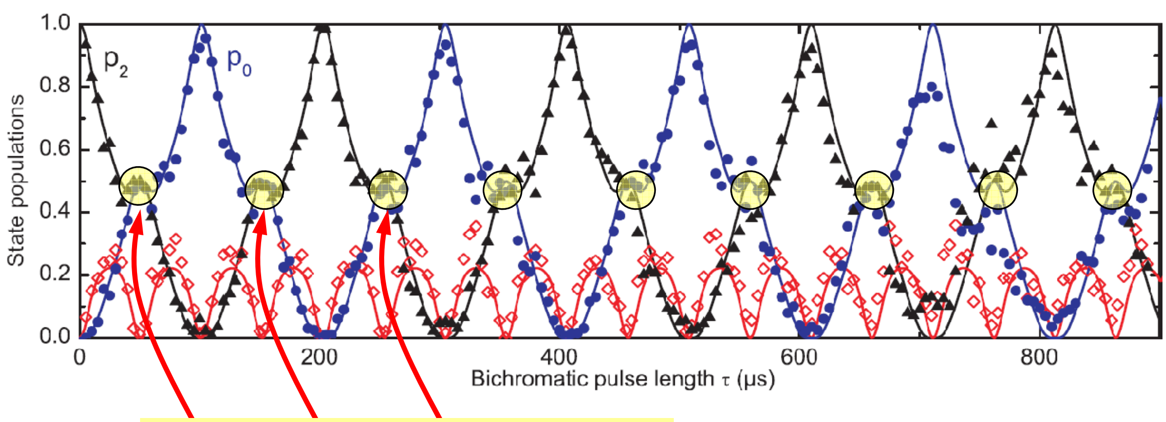
Theory: C. Roos, NJP **10,** 013002 (2008)

$$|00\rangle + |11\rangle \xrightarrow{R_2^C(\pi/2,\varphi),R_1^C(\pi/2,\varphi)}$$

$$\begin{aligned}
&\left(|0\rangle + ie^{i\varphi}|1\rangle\right)\left(|0\rangle + ie^{i\varphi}|1\rangle\right) \;+\; \left(|1\rangle + ie^{-i\varphi}|0\rangle\right)\left(|1\rangle + ie^{-i\varphi}|0\rangle\right) \\
= &\left(1 - e^{-2i\varphi}\right)|00\rangle + ie^{i\varphi}\left(1 + e^{-2i\varphi}\right)|01\rangle \\
&+ ie^{i\varphi}\left(1 + e^{-2i\varphi}\right)|10\rangle + \left(1 - e^{-2i\varphi}\right)|11\rangle\,,
\end{aligned}$$

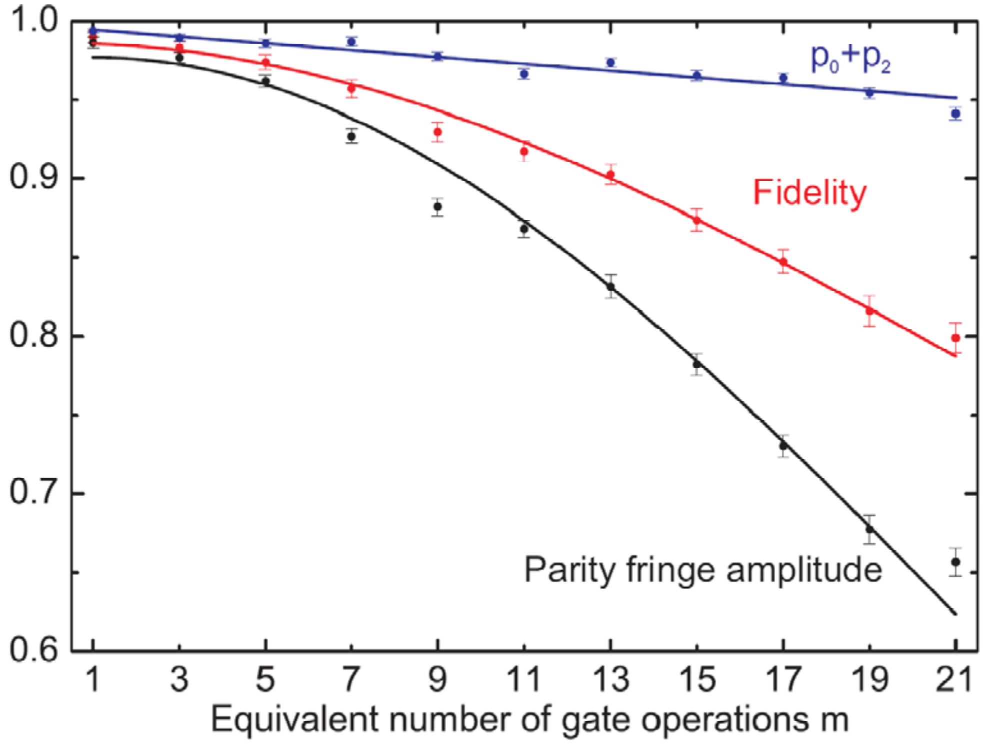


Gate concatenation



maximally entangled states

Gate performance



Scaling of this approach?

Problems :

- Coupling strength between internal and motional states of a N-ion string decreases as

$$\eta \propto \frac{1}{\sqrt{N}}$$

(momentum transfer from photon to ion string becomes more difficult)

-> Gate operation speed slows down

- More vibrational modes increase risk of spurious excitation of unwanted modes
- Distance between neighbouring ions decreases -> addressing more difficult

-> Use flexible trap potentials to split long ion string into smaller segments and perform operations on these smaller strings

DiVincenzo criteria

I. Scalable physical system, well characterized qubits ✓/ ?

II. Ability to initialize the state of the qubits ✓

III. Long relevant coherence times, much longer than gate operation time ✓

IV. "Universal" set of quantum gates ✓

V. Qubit-specific measurement capability ✓

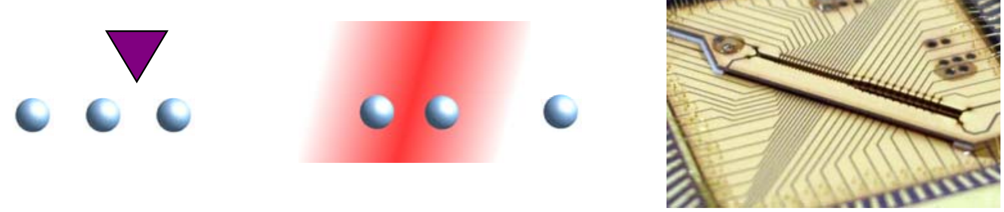
Often neglected:

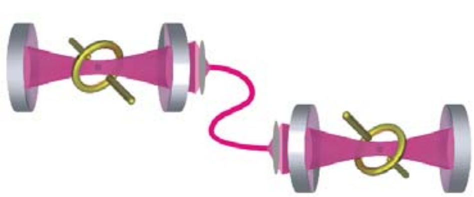
- exceptional fidelity of operations
- low error rate also for large quantum systems
- all requirements have to met at the same time

Scaling of ion trap quantum computers

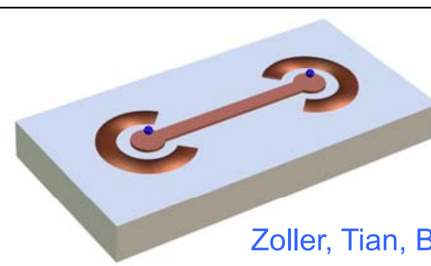
Its easy to have thousands of coherent qubits …
but hard to control their interaction

Kielpinski, Monroe, Wineland



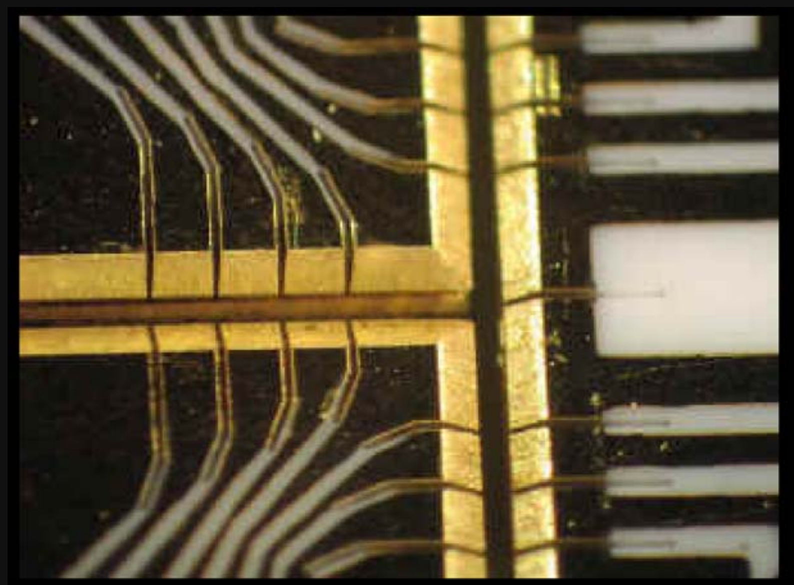


Cirac, Zoller, Kimble, Mabuchi



Zoller, Tian, Blatt

The Michigan T trap



An implementation of the Deutsch-algoritm …

Deutsch‘s problem: Introduction

Decide which class the coin is:
False (equal sides) or ***Fair***

Front

Back

A single measurement does NOT give the right answer

Deutsch‘s problem: Mathematical formulation

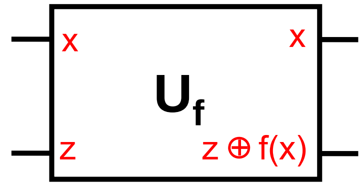
4 possible coins are representend by 4 functions

	Constant		**Balanced**	
	Case 1	**Case 2**	**Case 3**	**Case 4**
***f*(0)**	**0**	**1**	**0**	**1**
***f*(1)**	**0**	**1**	**1**	**0**

Deutsch‘s problem: Mathematical formulation

4 possible coins are representend by 4 functions

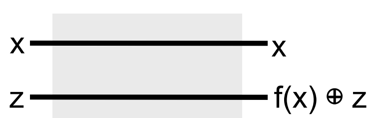
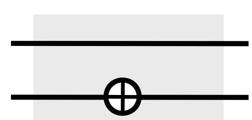
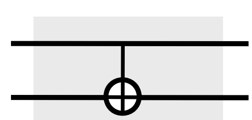
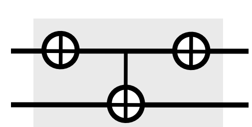
	Constant		Bal	
	Case 1	Case 2	Case 3	
f(0)	0	1	0	1
f(1)	0	1	1	0
z ⊕ *f*(*x*)	ID	NOT	CNOT	Z-CNOT



$$U_{fn}\left|x,z\right\rangle = \left|x, f_n(x) \oplus z\right\rangle$$

Physically reversible process
realized by a unitary transformation

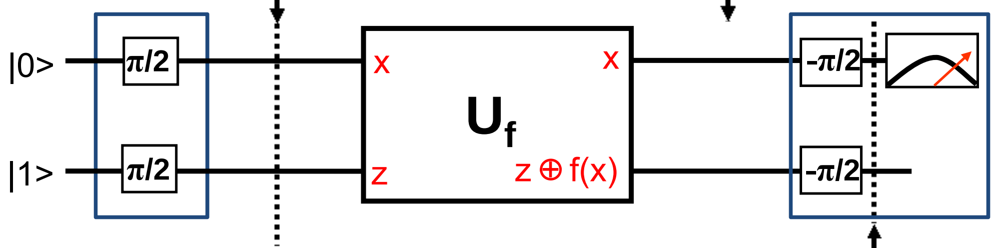
Deutsch Jozsa quantum circuit

Case	Logic	Quantum circuit	Matrix U_{fn}
f_1	*ID*		1000 0100 0010 0001
f_2	*NOT*		0100 1000 0001 0010
f_3	*CNOT*		1000 0100 0001 0010
f_4	*Z-CNOT*		0100 1000 0010 0001

Deutsch Jozsa quantum circuit

$|\psi_1\rangle = |0+1\rangle|0-1\rangle$

$|\psi_2\rangle = \pm|0+1\rangle|0-1\rangle$ if f(0) = f(1)
$= \pm|0-1\rangle|0-1\rangle$ if f(0) ≠ f(1)

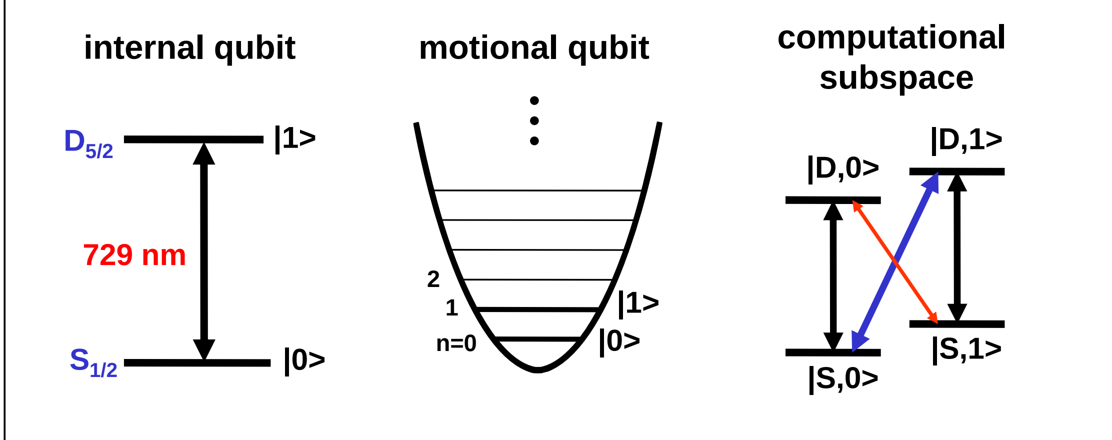


$|\psi_3\rangle = \pm|0\rangle|1\rangle$ if f(0) = f(1) *(Constant)*
$= \pm|1\rangle|1\rangle$ if f(0) ≠ f(1) *(Balanced)*

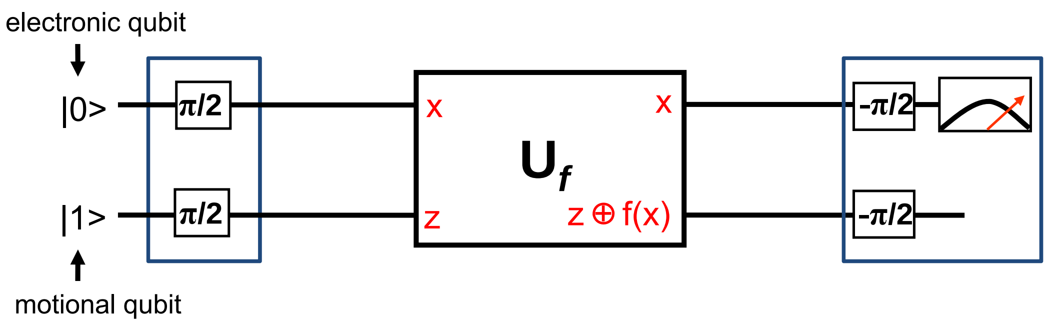
Quantum analysis gives the right answer after a single measurement!

- *D. Deutsch, R. Josza, Proc. R. Soc. London A439, 553 (1992)*
- *M. Nielsen, I. Chuang, QC and QI, Cambridge (2000)*

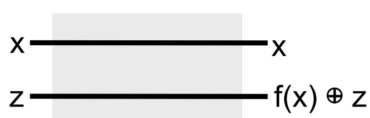
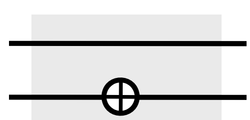
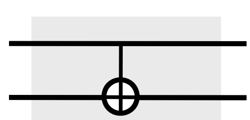
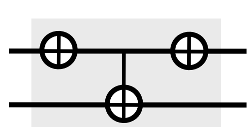
Qubits in $^{40}Ca^+$



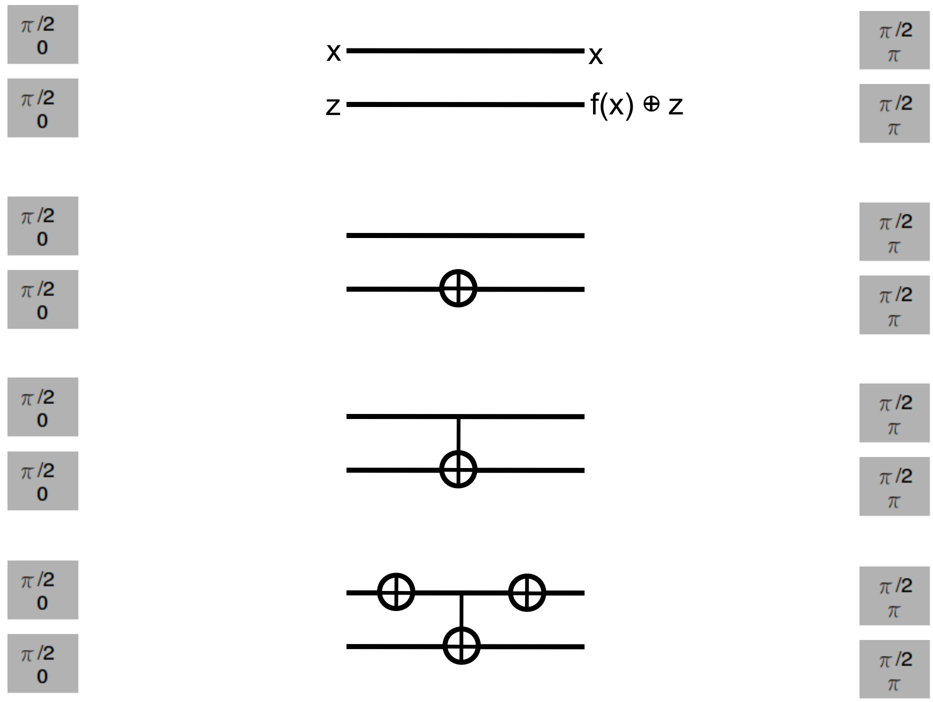
No information in the second qubit



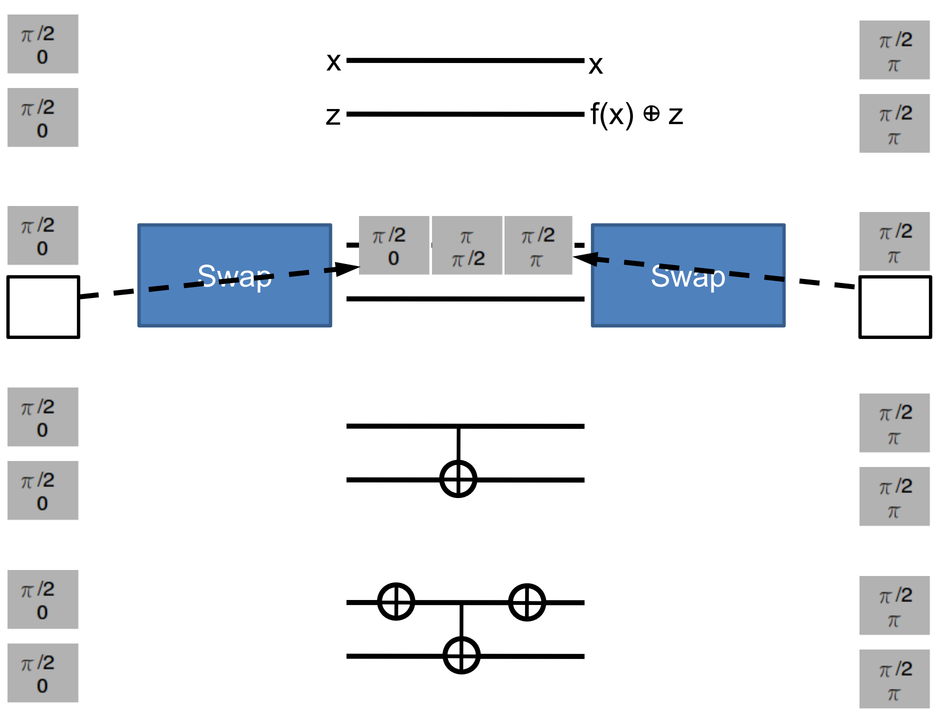
Deutsch Jozsa quantum circuit

Case	Logic	Quantum circuit	Matrix U_{fn}
f_1	*ID*		1000 0100 0010 0001
f_2	*NOT*		0100 1000 0001 0010
f_3	*CNOT*		1000 0100 0001 0010
f_4	*Z-CNOT*		0100 1000 0010 0001

Deutsch Jozsa: Realization

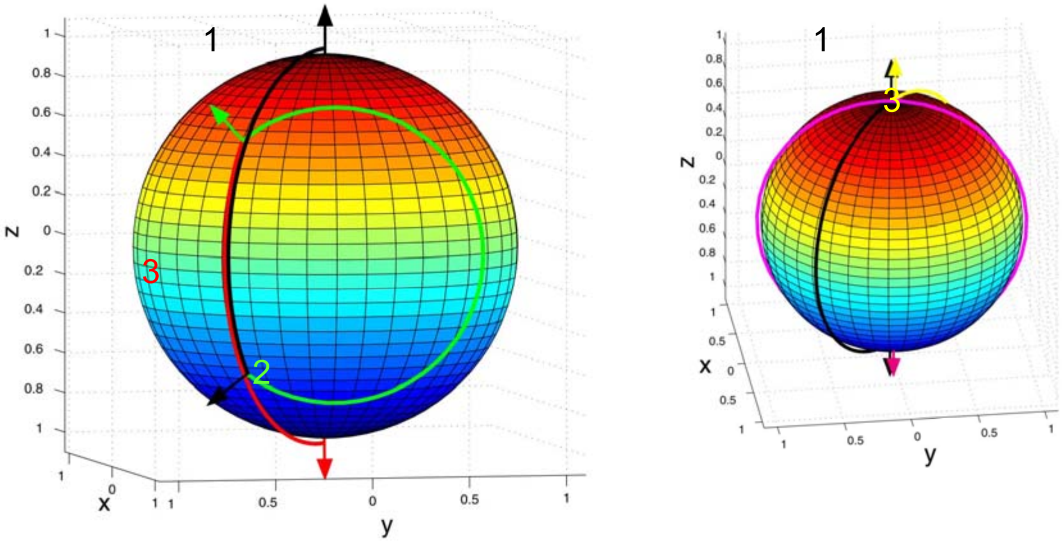


Deutsch Jozsa: Realization



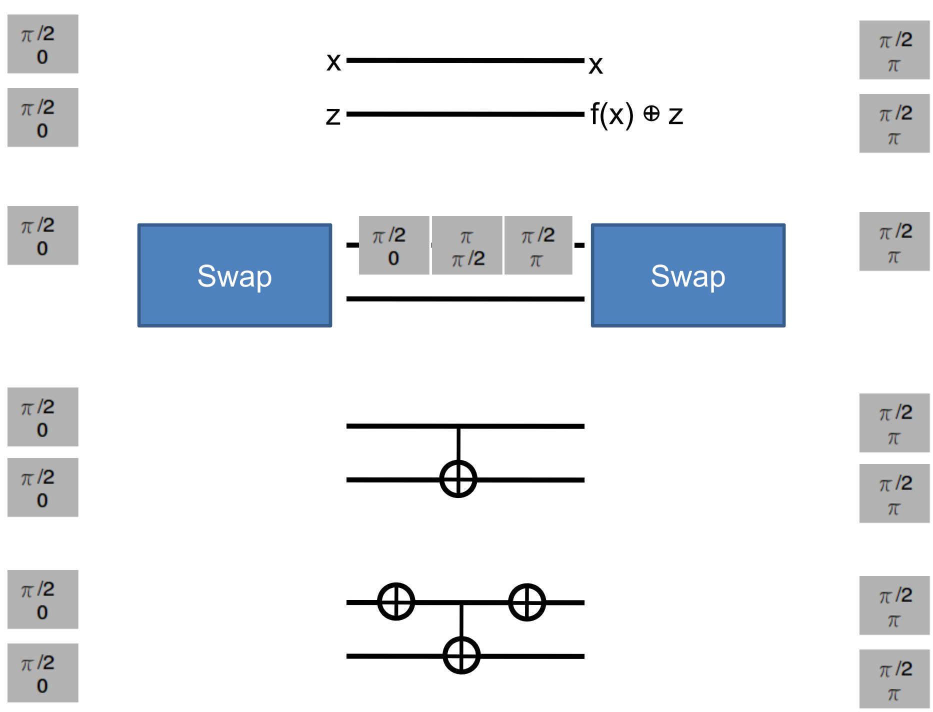
3-step composite SWAP operation

$$R^+\left(\frac{\pi}{\sqrt{2}},\pi\right) R^+\left(\frac{2\pi}{\sqrt{2}},\pi+\varphi_{\text{swap}}\right) R^+\left(\frac{\pi}{\sqrt{2}},\pi\right)$$

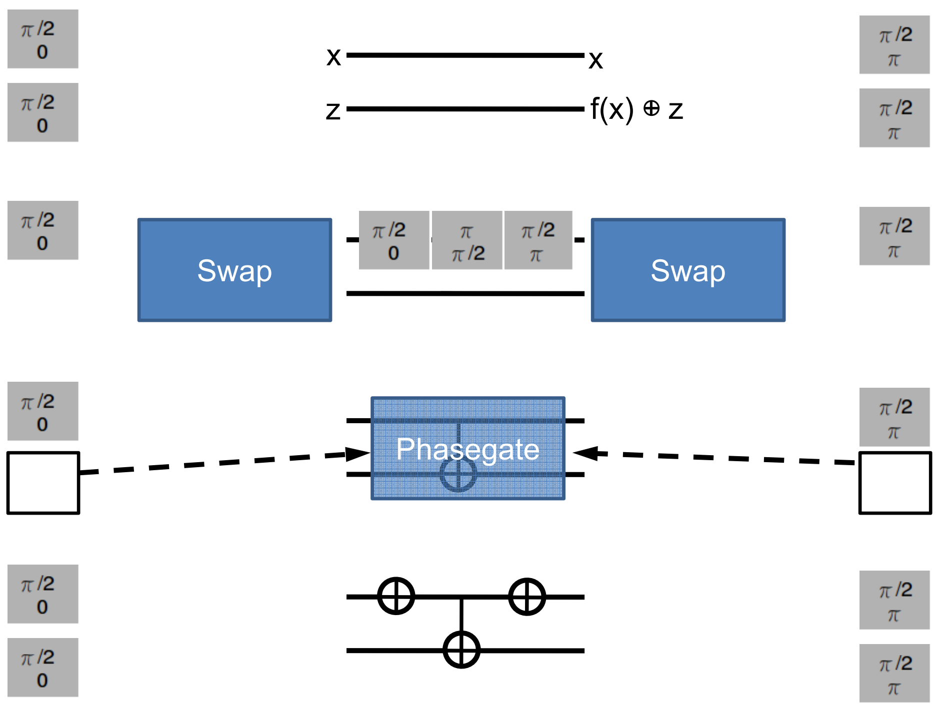


I. Chuang et al., Innsbruck (2002)

Deutsch Jozsa: Realization

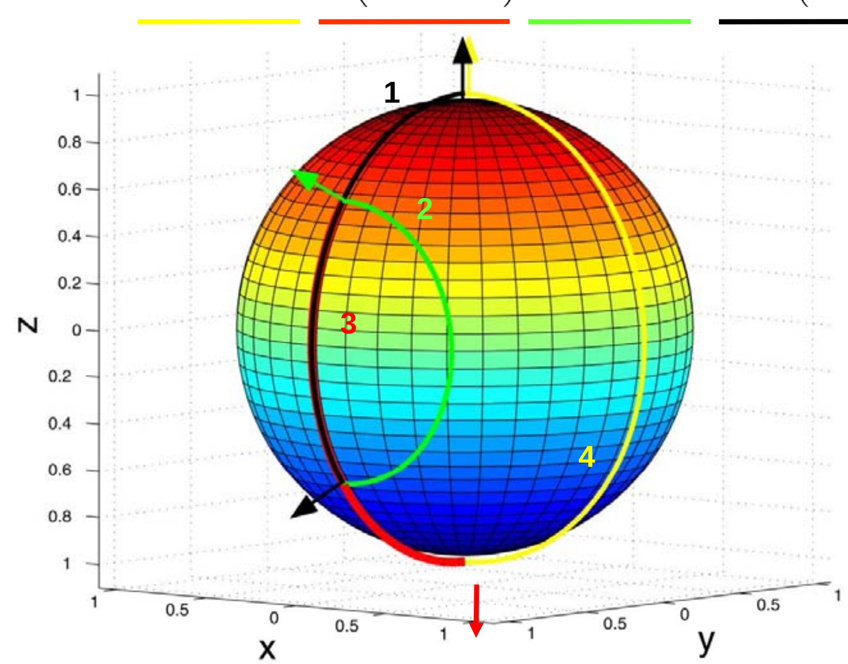


Deutsch Jozsa: Realization



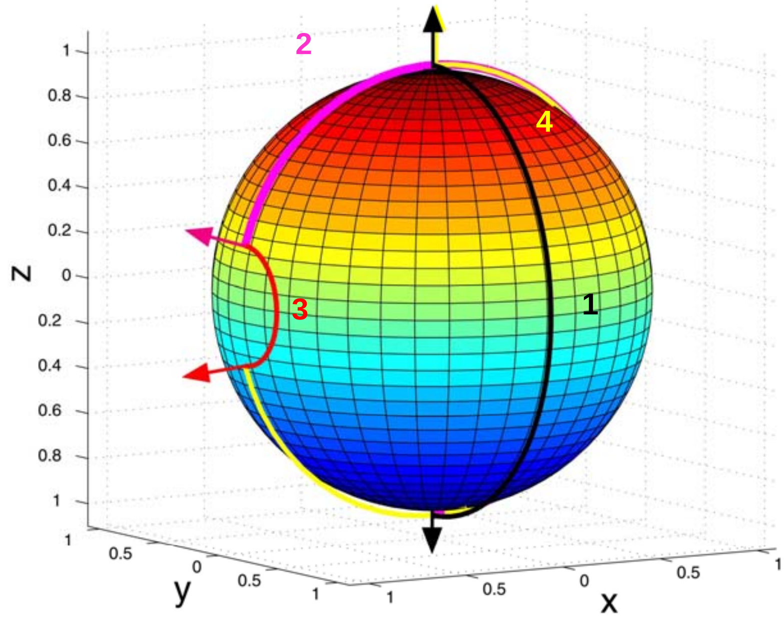
Composite phase gate (2π rotation)

$$R^{+}(\pi,\pi/2)\,R^{+}\left(\pi/\sqrt{2},0\right)\,R^{+}(\pi,\pi/2)\,R^{+}\left(\pi/\sqrt{2},0\right)$$

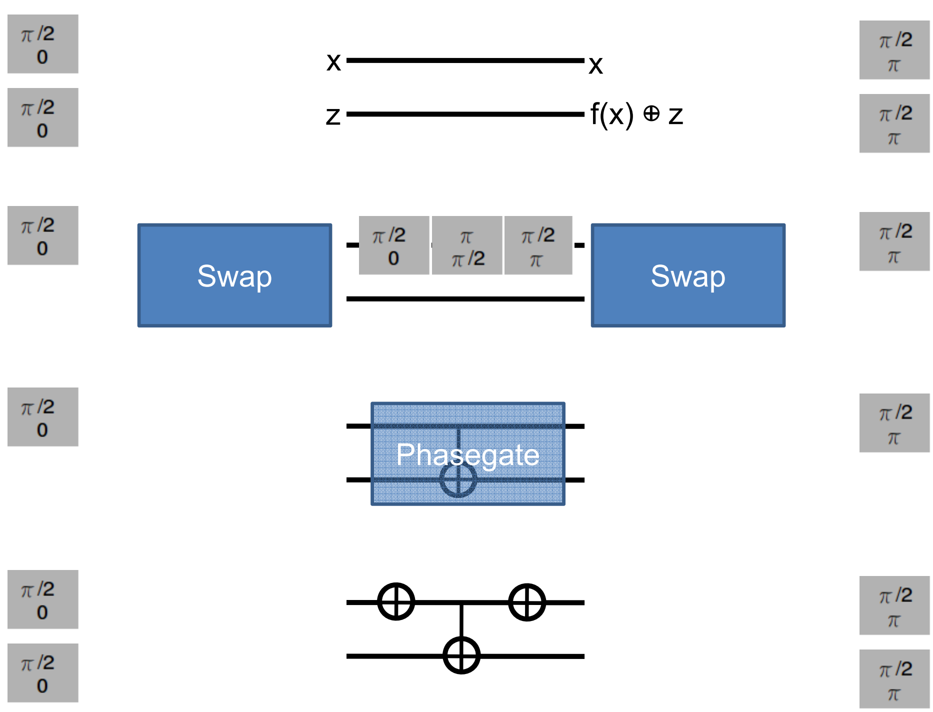


Action on |S,1> - |D,2>

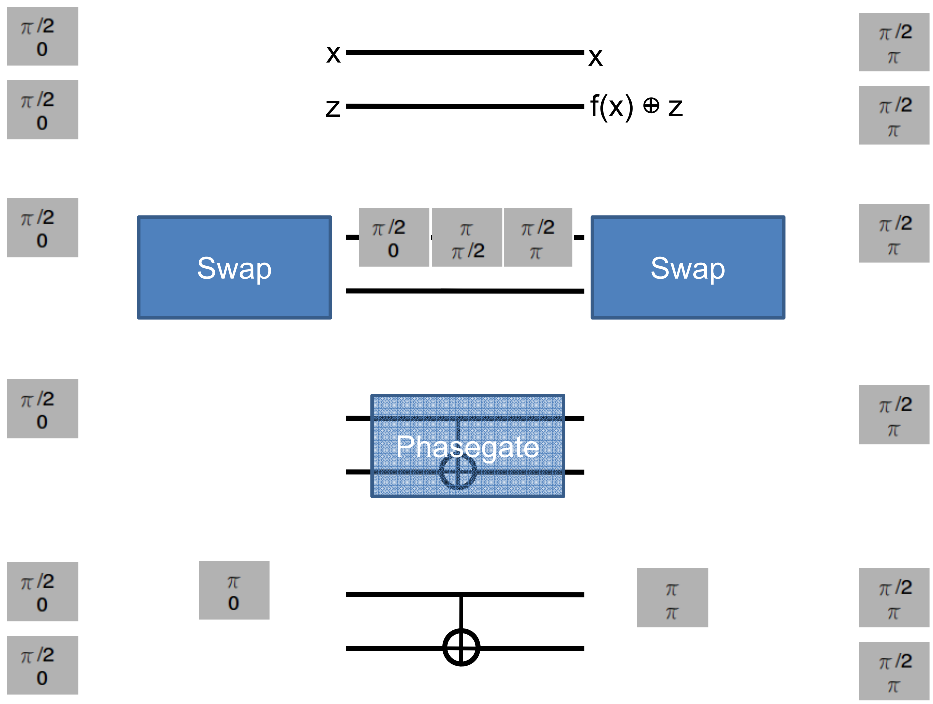
$$R^{+}(\pi,\pi/2)\,R^{+}\left(\pi/\sqrt{2},0\right)\,R^{+}(\pi,\pi/2)\,R^{+}\left(\pi/\sqrt{2},0\right)$$



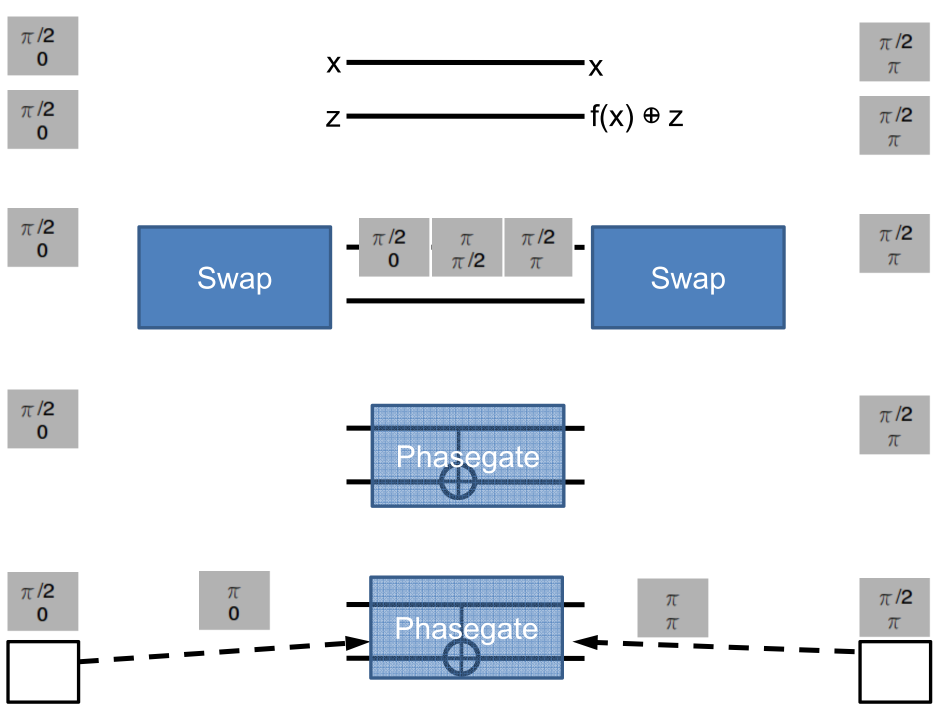
Deutsch Jozsa: Realization



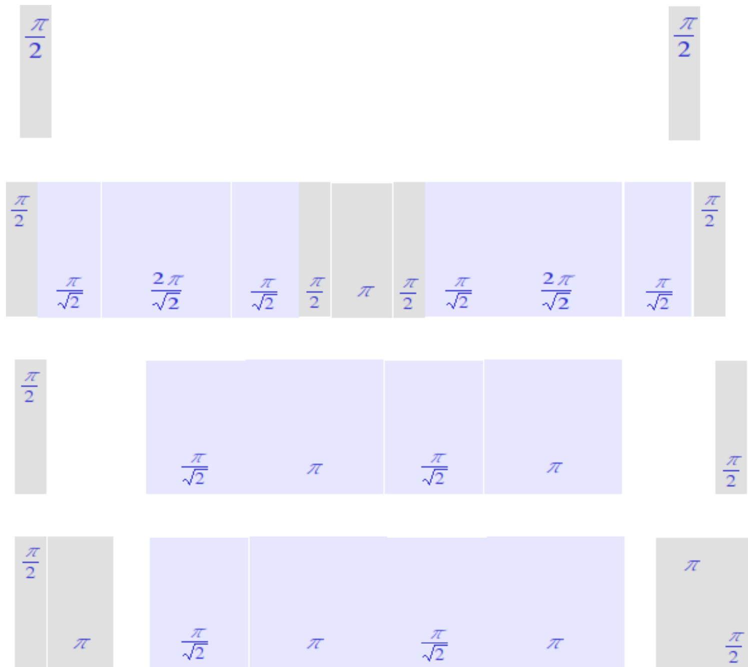
Deutsch Jozsa: Realization



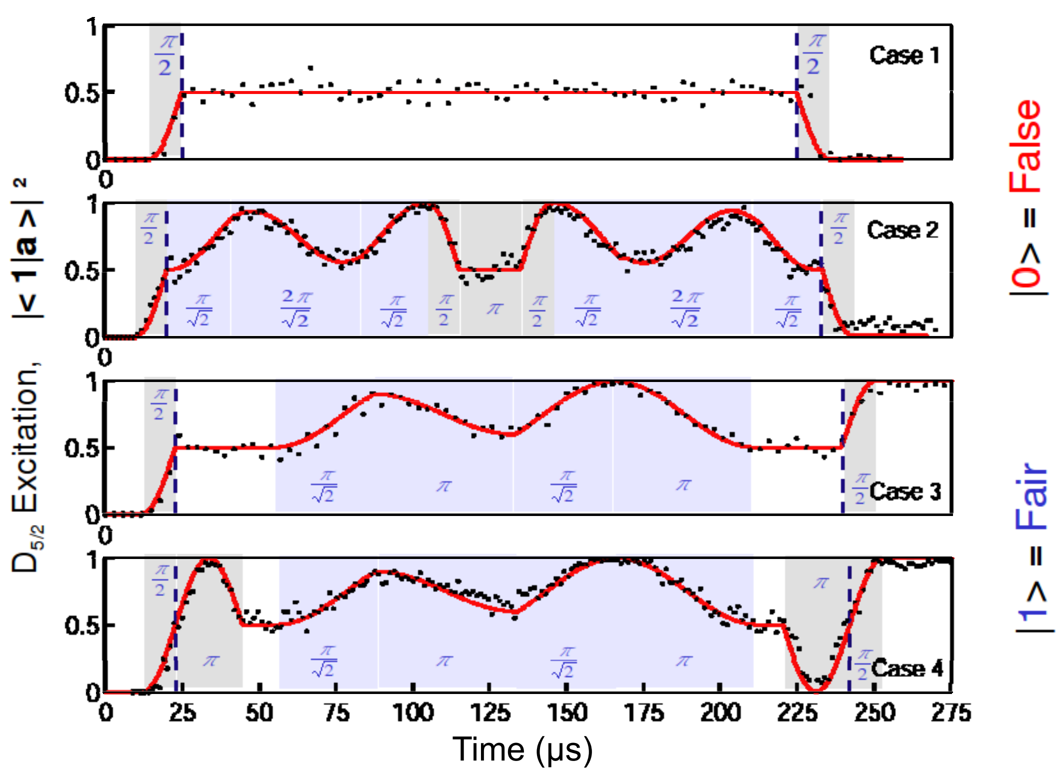
Deutsch Jozsa: Realization



Deutsch Jozsa: Realization



Deutsch Jozsa: Realization



Deutsch Jozsa: Result

	Constant		Balanced	
	Case 1	Case 2	Case 3	Case 4
expected $\lvert\langle 1 \vert a\rangle\rvert^2$	0	0	1	1
measured $\lvert\langle 1 \vert a\rangle\rvert^2$	0.019(6)	0.087(6)	0.975(4)	0.975(2)
expected $\lvert\langle 1 \vert w\rangle\rvert^2$	1	1	1	1
measured $\lvert\langle 1 \vert w\rangle\rvert^2$	--	0.90(1)	0.931(9)	0.986(4)

S. Gulde et al., Nature 412, 48 (2003)

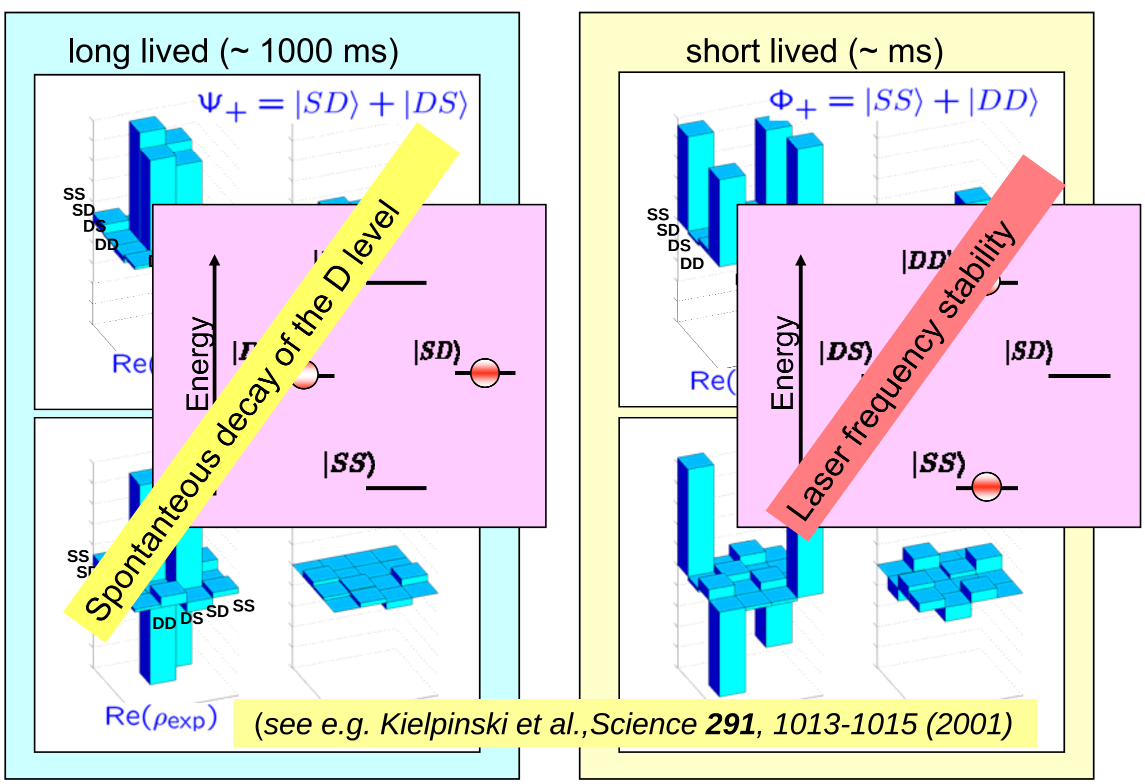
Conclusions

- Basics of ion trap quantum computing
- Measuring a density matrix
- Quantum gates
- Deutsch Algorithm

€ FWF SFB
SCALA QGATES
Industrie Tirol
IQI GmbH
FWF bm:bwk
$ IARPA

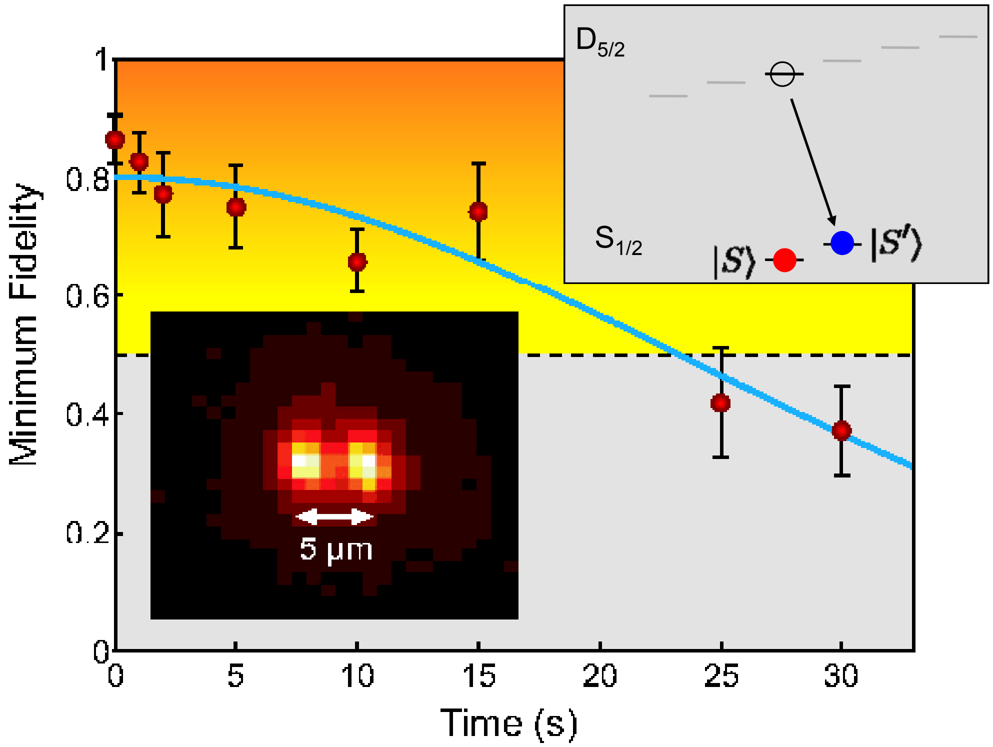
Berkeley, Nov 25th 2008

Decoherence properties of qubits



(see e.g. Kielpinski et al.,Science **291**, 1013-1015 (2001))

Ultra long lived entanglement



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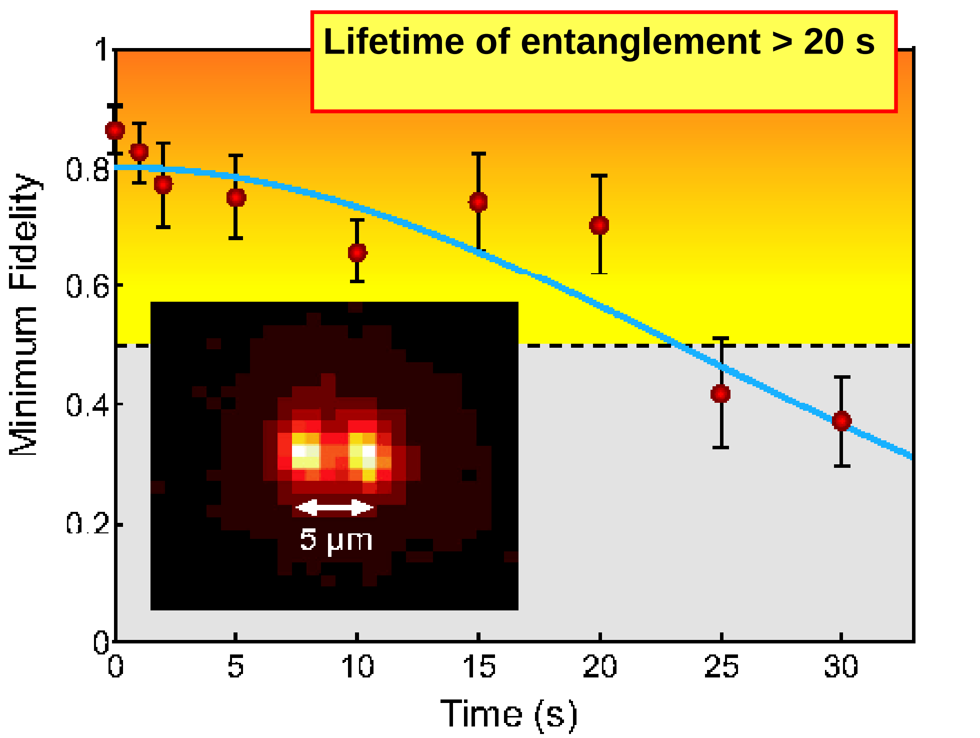
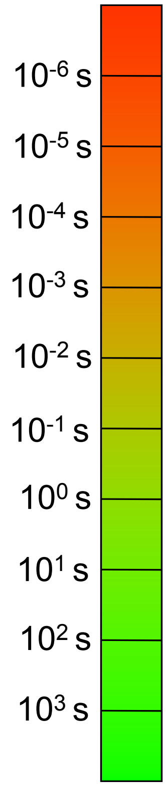


H. Häffner et al., App. Phys. B **81** 151 (2005).

Realized time scales



Single qubit gates

Two qubit gates (Geometric phase gates: Boulder, Innsbruck)

Two qubit gates (Cirac-Zoller approach)

Single qubit coherence (magnetic field sensitive)

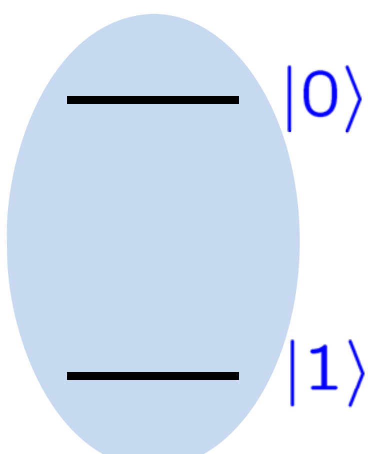
Coherence of the motion (Innsbruck)

Single qubit coherence (magnetic field insensitive, Boulder)
Decoherence free subspaces

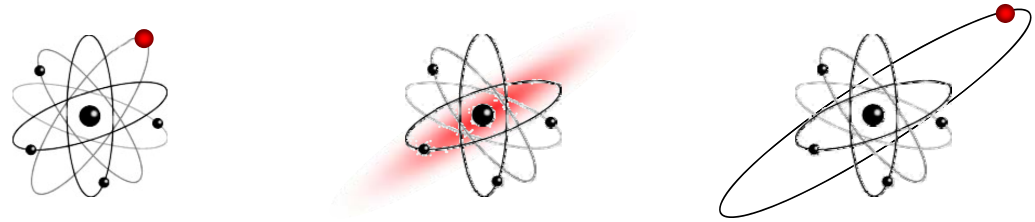
Well chosen single qubit coherence (Boulder)

A quantum bit

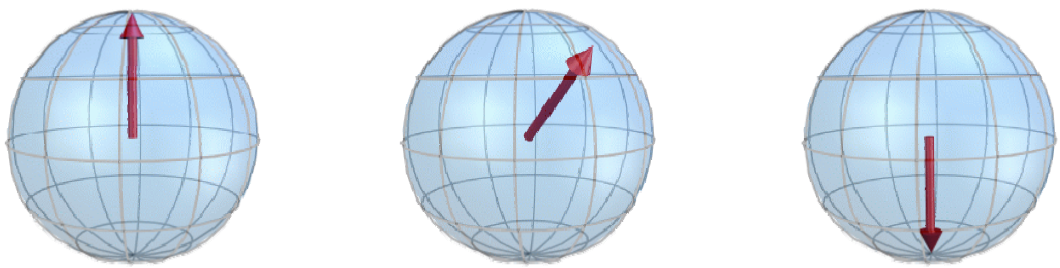
Two level system:



Qubits



$|0\rangle$ $\qquad \alpha|0\rangle + \beta|1\rangle$ $\qquad |1\rangle$



Possible qubit encodings

Physical Qubit $\longleftrightarrow$ **Logical Qubit**

$|0\rangle_P = |D\rangle$ $\qquad |0\rangle_L = |SD\rangle$

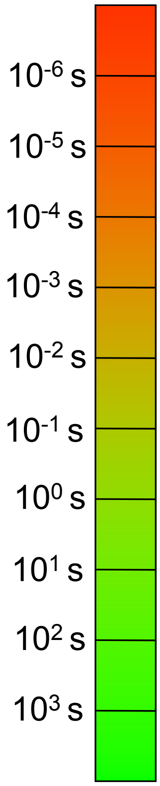
$|1\rangle_P = |S\rangle$ $\qquad |1\rangle_L = |DS\rangle$

Effect of magnetic field or laser frequency fluctuations on qubits

$$\begin{array}{c} |D\rangle + |S\rangle \\ \Downarrow \\ e^{i\phi}|D\rangle + |S\rangle \end{array} \quad \longleftrightarrow \quad \begin{array}{c} |SD\rangle + |DS\rangle \\ \Downarrow \\ e^{i\phi}\left(|SD\rangle + |DS\rangle\right) \end{array}$$

Logical qubit experiences global phase only

Realized time scales



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Two qubit gates (Cirac-Zoller approach)

Single qubit coherence (magnetic field sensitive)

Coherence of the motion (Innsbruck)

Single qubit coherence (magnetic field insensitive)

Decoherence free subspaces

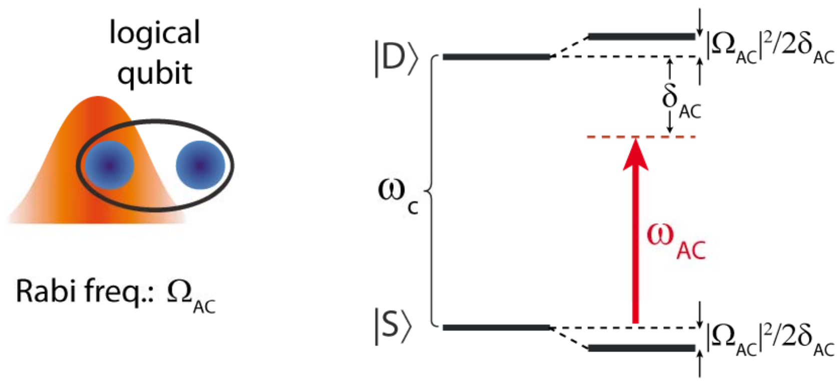
Well chosen single qubit (Boulder)

Universal set of gates in a DFS

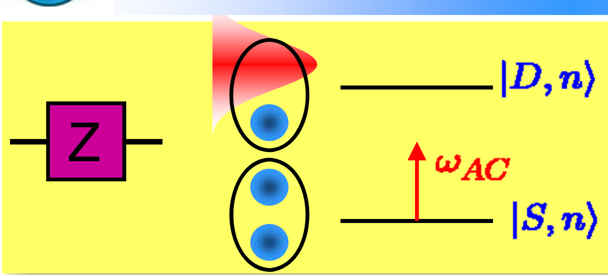
- single qubit operations
  - Z gates
  - X gates

- two –qubit operations
  - phase gate

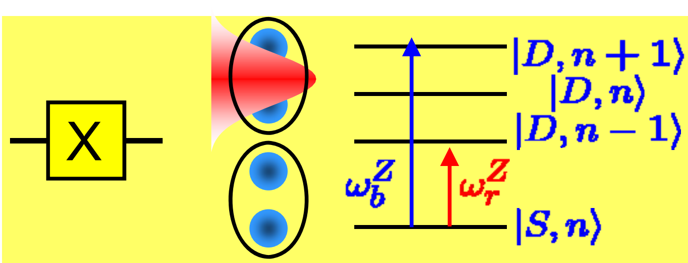
Z gate



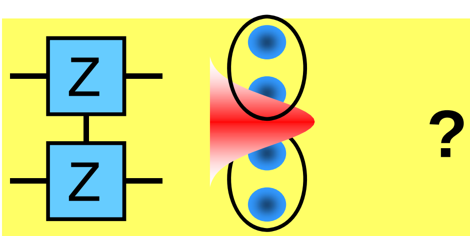
Universal set of gates in a DFS



single qubit phaseshift by
AC Stark shifts

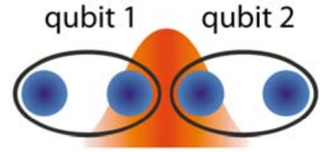


$$\omega_b^{MS} = \omega_c + \nu + \delta_{MS}$$
$$\omega_r^{MS} = \omega_c - \nu - \delta_{MS}$$



Two logical qubit gate

Two body interactions preferred:



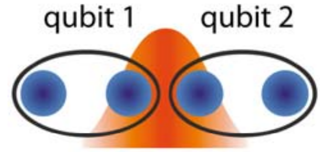
Most interactions cause the state to leave the decoherence free subspace.

Some solutions: L. Aolita et al., PRA **75** 052337 (2007)

Two logical qubit phase gate

Action of the phase gate on two physical qubits:

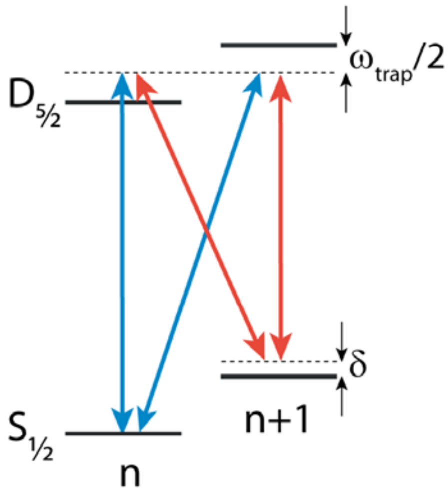
$$\begin{array}{c} |DD\rangle \\ |DS\rangle \\ |SD\rangle \\ |SS\rangle \end{array} \Rightarrow \begin{array}{c} e^{i\phi}|DD\rangle \\ |DS\rangle \\ |SD\rangle \\ e^{i\phi}|SS\rangle \end{array}$$

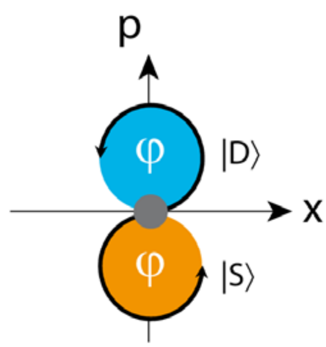


...and on the logical qubits:

$$\begin{array}{c} |00\rangle_l \\ |01\rangle_l \\ |10\rangle_l \\ |11\rangle_l \end{array} = \begin{array}{c} |S\rangle|DS\rangle|D\rangle \\ |S\rangle|DD\rangle|S\rangle \\ |D\rangle|SS\rangle|D\rangle \\ |D\rangle|SD\rangle|S\rangle \end{array} \Rightarrow \begin{array}{c} |S\rangle|DS\rangle|D\rangle \\ |S\rangle e^{i\phi}|DD\rangle|S\rangle \\ |D\rangle e^{i\phi}|SS\rangle|D\rangle \\ |D\rangle|SD\rangle|S\rangle \end{array} = \begin{array}{c} |00\rangle_l \\ e^{i\phi}|01\rangle_l \\ e^{i\phi}|10\rangle_l \\ |11\rangle_l \end{array}$$

Two qubit phase gate

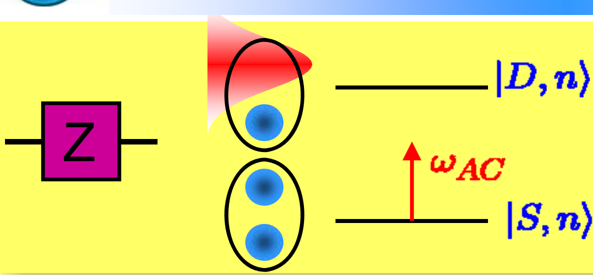




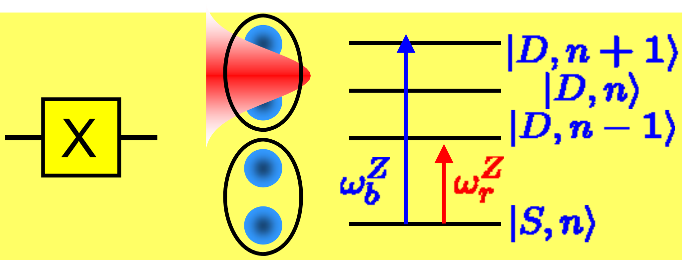
D. Leibfried, et al., Nature **422** 412 (2003)

K. Kim et. al., Phys. Rev. A**77**, 050303 (2008)

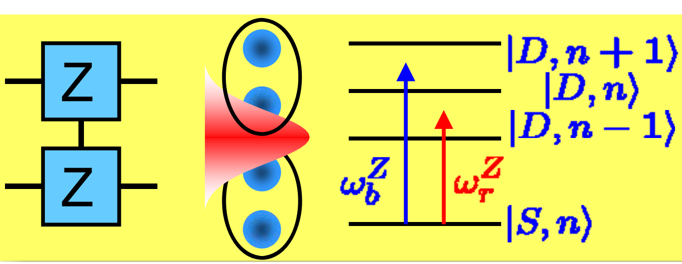
Universal set of gates in a DFS



single qubit phaseshift by AC Stark shifts

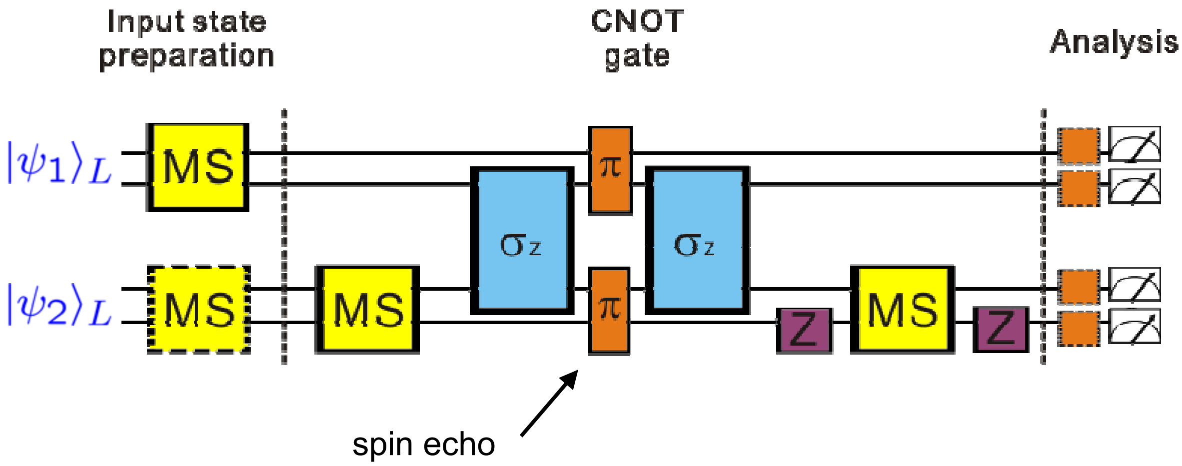


$$\omega_b^{MS} = \omega_c + \nu + \delta_{MS}$$
$$\omega_r^{MS} = \omega_c - \nu - \delta_{MS}$$

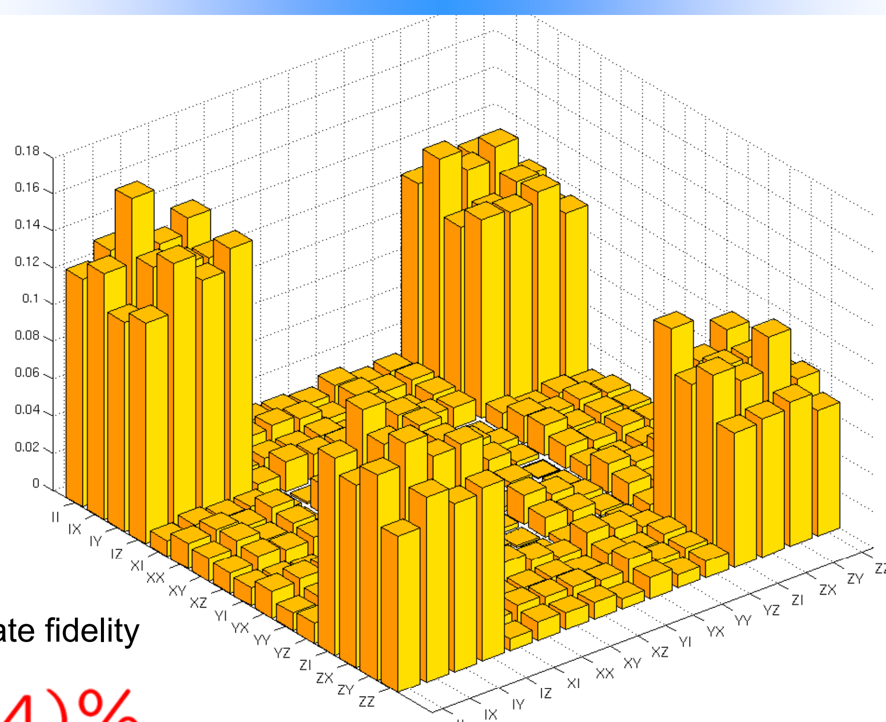


$$\omega_b^Z = \omega_c + \nu/2 + \delta_z$$
$$\omega_r^Z = \omega_c - \nu/2 - \delta_z$$

A CNOT in a DFS



Process tomography of the CNOT



mean gate fidelity

89(4)%

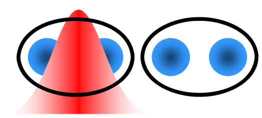
Discussion

mean gate fidelity: 89(4)%

(after DFS postselection)

Main limitations:

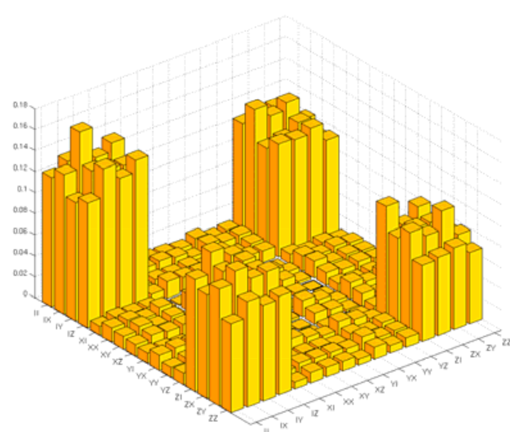
- spurious laser frequency components
- off-resonant coupling to other levels
- intensity stability on ions
- addressing errors



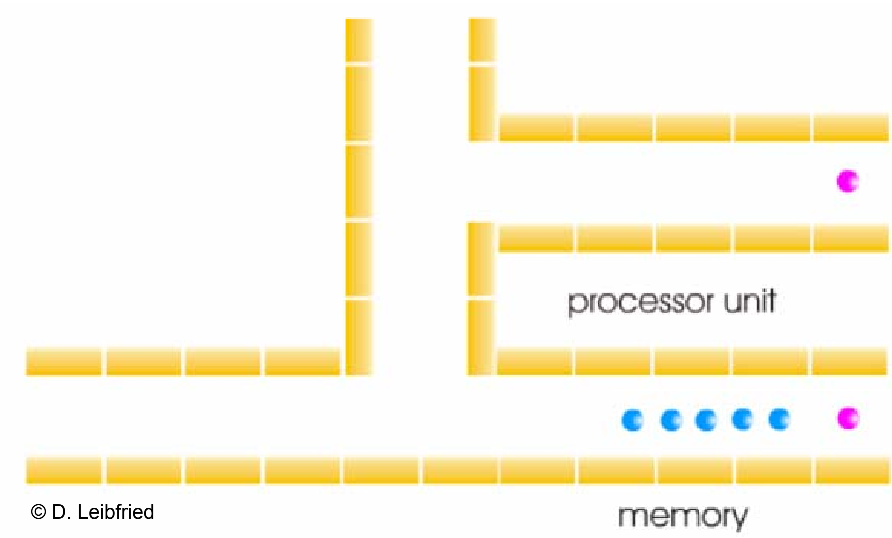
Discussion

Advantages:

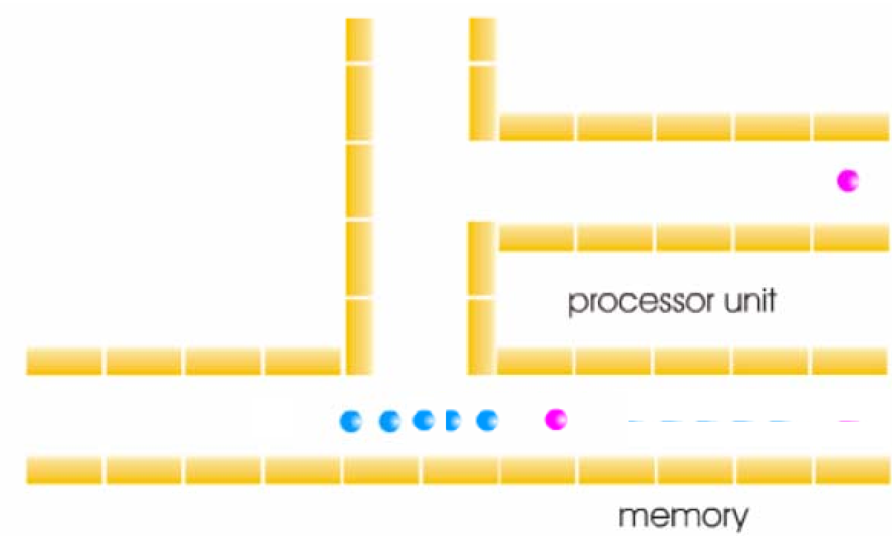
- lifetime limited coherence time
- insensitive to laser linewidth
- insensitive to AC-Stark shifts



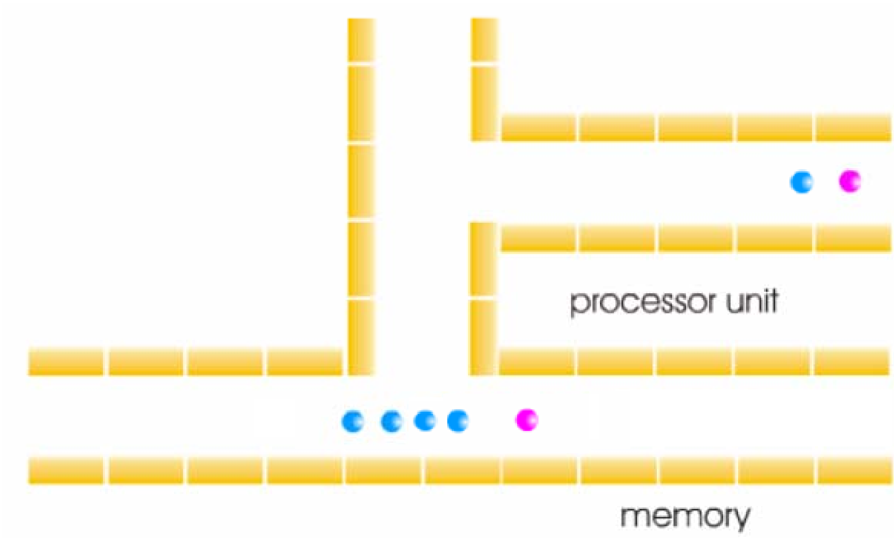
Scaling of ion trap quantum computers



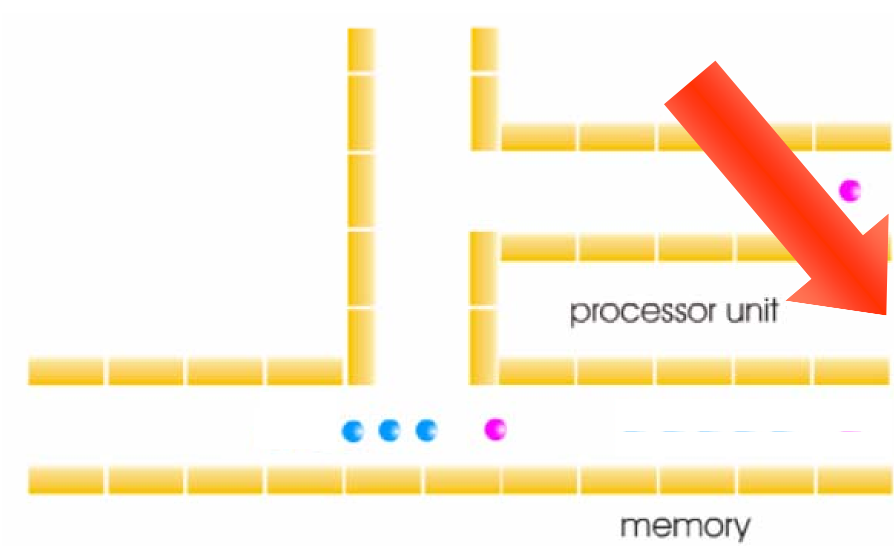
Scaling of ion trap quantum computers



Scaling of ion trap quantum computers

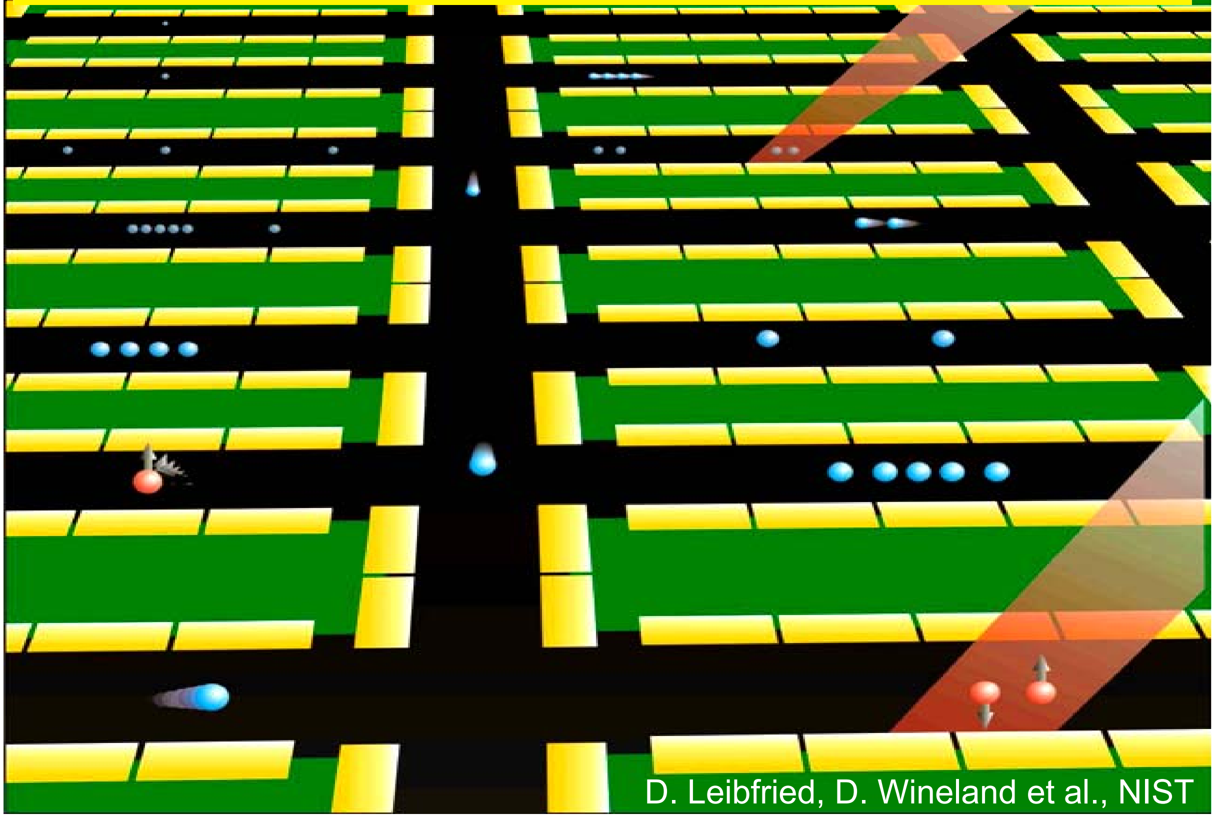


Scaling of ion trap quantum computers



„Architecture for a large-scale ion-trap quantum computer“, D. Kielpinski et al, Nature **417**, 709 (2002)

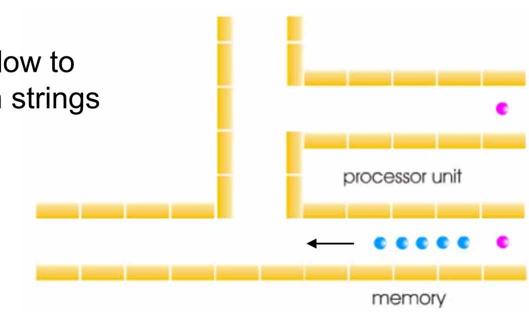
Multiplexed trap structure: NIST Boulder



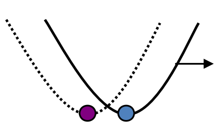
Segmented ion traps as scalable trap architecture

(ideas pioneered by D. Wineland, NIST)

Segmented trap electrode allow to transport ions and to split ion strings



State of the art:



Transport of ions

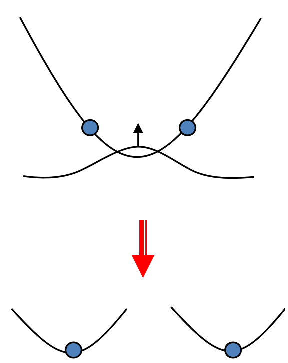
1 mm within 50 μs

no motional heating

Splitting of two-ion crystal

$t_{separation} \approx 200\ \mu s$

small heating $n \approx 1$



„Architecture for a large-scale ion-trap quantum computer“, D. Kielpinski et al, Nature **417**, 709 (2002)

„Transport of quantum states“, M. Rowe et al, quant-ph/0205084

Scaling of this approach?

Problems :

- Coupling strength between internal and motional states of a N-ion string decreases as

$$\eta \propto \frac{1}{\sqrt{N}}$$

(momentum transfer from photon to ion string becomes more difficult)

-> Gate operation speed slows down

- More vibrational modes increase risk of spurious excitation of unwanted modes

- Distance between neighbouring ions decreases -> addressing more difficult

-> Use flexible trap potentials to split long ion string into smaller segments and perform operations on these smaller strings